

## A topological optimization method for flexible multi-body dynamic system using epsilon algorithm

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**Abstract.** In a flexible multi-body dynamic system the typical topological optimization method for structures cannot be directly applied, as the stiffness varies with position. In this paper, the topological optimization of the flexible multi-body dynamic system is converted into structural optimization using the equivalent static load method. First, the actual boundary conditions of the control system and the approximate stiffness curve of the mechanism are obtained from a flexible multi-body dynamical simulation. Second, the finite element models are built using the absolute nodal coordination for different positions according to the stiffness curve. For efficiency, the static reanalysis method is utilized to solve these finite element equilibrium equations. Specifically, the finite element equilibrium equations of key points in the stiffness curve are fully solved as the initial solution, and the following equilibrium equations are solved using a reanalysis method with an error controlled epsilon algorithm. In order to identify the efficiency of the elements, a non-dimensional measurement is introduced. Finally, an improved evolutionary structural optimization (ESO) method is used to solve the optimization problem. The presented method is applied to the optimal design of a die bonder. The numerical results show that the presented method is practical and efficient when optimizing the design of the mechanism.

**Keywords:** topological optimization; the flexible multi-body dynamic system; die bonder; epsilon algorithm.

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### 1. Introduction

The structural optimization problem has been a subject under numerous studies in recent years. Topological optimization can greatly improve the design and potential savings are generally more significant than those resulting from fixed-topology optimization (Liu and Chen 1992). It can also be used for optimal design of geometrically and materially nonlinear structures (Huang and Xie 2010), and dynamic problem with non-stochastic structural uncertainty (Lee *et al.* 2010). However, the topological optimization method is difficult to be applied to dynamic response optimization, one of the main difficulties is the design sensitivity analysis, which is inevitable in gradient-based optimization methodology. Further difficulties involve the treatment of the dynamic constraints, and the computational solution during the optimization process.

To evaluate design sensitivity in a dynamic response optimization, one has to solve many

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differential equations resulting from the number of design variables and the number of active constraints (Arora 1999). If a transformation method such as the augmented Lagrangian method (ALM) is used with the adjoint variable method (AVM), the cost of calculating the design sensitivity can be reduced. However, a terminal value problem must be solved to acquire the adjoint variable for the augmented functional. If the terminal value has some error that has been accumulating over the entire time interval where the equations of motion are integrated, the adjoint variable for the augmented functional can diverge or be inaccurate. Another drawback of the ALM in dynamic response optimization is that some multiplier updating rules require the gradient of an individual response. These drawbacks reduce the attraction the ALM and AVM methods, although they considerably reduced computational effort.

Another predicament we often encounter in structural dynamic response optimization is how to treat dynamic constraints. To remove the time parameter from the dynamic constraints, it is customary to replace the dynamic constraint with an equivalent integral form or several pointwise constraints (Tseng and Arora 1989). In general, this replacement causes an increase in the number of constraints or numerical instability in the optimization procedure. We need extra effort to reduce the number of constraints or to stabilize the optimization procedure. There are many optimization examples of large structures that are subject to static loads. The existence of many practical examples means that the static response optimization methodology is well established. Design sensitivity evaluation in static response optimization does not need to solve differential equations but only simpler simultaneous algebraic equations. In the modern computing environment, it is relatively easier to solve large-scale algebraic equations. If the advantages of static response optimization can be exploited in dynamic response optimization situations, it will be much easier to perform dynamic response optimization of a structure. Basically the objective of the treatment of a dynamic constraint is the elimination of the time parameter. After eliminating the time parameter, we try to find the optimum design by using various optimization techniques. At this point, if we have some compensation for the errors of design sensitivity and response, the static response optimization process can be adopted in dynamic response optimization situations. Recently, a structural dynamic response optimization procedure using static response optimization techniques has been demonstrated (Kang *et al.* 2001, Choi and Park 2002). The main idea is transformation of a dynamic load into equivalent static loads and performing a static response optimization with the transformed equivalent static load set. This method is applied to optimization of structure with dynamic loads (Park and Kang 2003) and flexible multi-body dynamic systems (Kang *et al.* 2005). Optimization of large-scale structures using conventional formulations often involves much computational effort. Repeated solutions of the analysis and sensitivity analysis equations usually require most of this effort. The computational cost may become prohibitive in large-scale structures having complex analysis models. To alleviate this difficulty, various procedures are developed and integrated into a general optimization approach. The approach is suitable for different classes of response types and optimization methods, including linear and non-linear response (Kirsch and Bogomolni 2007, Bogomolni *et al.* 2006); static and dynamic response (Kirsch *et al.* 2006, 2007); and direct and gradient optimization methods (Kirsch *et al.* 2005). Combined approximations are used for reanalysis and repeated sensitivity analysis (Kirsch and Bogomolni 2004). The advantage is that the efficiency of local approximations and the improved quality of global approximations are combined to obtain effective solution procedures. Approximate reanalysis and finite-difference sensitivity reanalysis are considered for each intermediate design during the solution process. Reductions in the computational effort may be several orders of magnitude less than would

otherwise be required. The Epsilon algorithm can accelerate the convergence speed, need less computer effort and can obtain higher accuracy even when the changes are large (Wu *et al.* 2007, Chen *et al.* 2006).

In this paper, the problem of the mechanism's structural optimization is treated as that of structural optimization using the equivalent static load method. Firstly, the actual boundary conditions of the control system and the approximate stiffness curve of the mechanism are obtained from a flexible multi-body dynamical simulation. Then according to the stiffness curve, the governing equations of the mechanism are divided into separate finite element equilibrium equations for different positions. For efficiency, the static reanalysis method is utilized to solve these equations. The detailed procedure is that, the finite element equilibrium equations of the key points in the stiffness curve are fully solved as the initial solution, and the following equilibriums are solved using a reanalysis method with the Epsilon algorithm. In order to identify the efficiency of elements, a non-dimensional measurement is introduced, and an improved ESO method is used to solve the optimization problem. The presented method is applied to the optimal design for a die bonder. The numerical results show that the method is practicable and efficient.

## 2. Technical background

For mechanisms operating at very high speed, the vibration of the structures must be considered, when the deformation is small, the relative nodal coordinate formulation (RNCF) with a floating frame of reference extended from the motion of the rigid body system is used Eq. (1), while for large deformations the absolute nodal coordinate formulation (ANCF) turned out to be very efficient Eq. (4).

According to previous work (Kubler 2003), the equation for the relative nodal coordinate formulation extended from the motion of the rigid body system is

$$\mathbf{M}(\mathbf{r}, t)\ddot{\mathbf{r}} + \mathbf{K}(\mathbf{r}, \dot{\mathbf{r}}, t) + \mathbf{K}_i(\mathbf{r}, \ddot{\mathbf{r}}) = \mathbf{q}(\mathbf{r}, \dot{\mathbf{r}}, t) \quad (1)$$

Where  $\mathbf{M}(\mathbf{r}, t)$  is the mass matrix of the rigid body system,  $\mathbf{K}(\mathbf{r}, \dot{\mathbf{r}}, t)$ ,  $\mathbf{K}_i(\mathbf{r}, \ddot{\mathbf{r}})$  and  $\mathbf{q}(\mathbf{r}, \dot{\mathbf{r}}, t)$  are the generalized elastic force, generalized Coriolis force and external force of the rigid body with position  $\mathbf{r}$ , respectively. The velocity and acceleration are  $\dot{\mathbf{r}}$  and  $\ddot{\mathbf{r}}$ , respectively.

In comparison to the equation of motion of the rigid body system the additional term

$$\mathbf{K}_i(\mathbf{r}, \ddot{\mathbf{r}}) = \begin{bmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{K} \end{bmatrix} \mathbf{r} + \begin{bmatrix} \mathbf{0} & \mathbf{0} \\ \mathbf{0} & \mathbf{D} \end{bmatrix} \dot{\mathbf{r}} \quad (2)$$

depends only on the stiffness and damping matrices  $\mathbf{K}$  and  $\mathbf{D}$  of the flexible bodies. Moreover, the inertial matrix shows the inertial coupling due to the relative coordinates

$$\mathbf{M} = \begin{bmatrix} \mathbf{M}_{rr} & \mathbf{M}_{rf} \\ \mathbf{M}_{rf}^T & \mathbf{M}_{ff} \end{bmatrix} \quad (3)$$

where lower case  $r$  and  $f$  denote the rigid bodies and flexible body, respectively. Within ANCF for highly flexible bodies absolute coordinates are summarized in a vector  $\mathbf{r}_a$  characterizing the material points of the bodies by an appropriate shape function. The equation of motion is

$$\mathbf{M}\ddot{\mathbf{r}}_a + \mathbf{K}_a(\mathbf{r}_a)\mathbf{r}_a = \mathbf{q}(\mathbf{r}_a, t) \quad (4)$$

where  $\mathbf{M}$  is a constant mass matrix and the vector  $\mathbf{K}_i(\mathbf{r}, \dot{\mathbf{r}})$  of the generalized Coriolis force vanishes due to the absolute coordinates. The equation of motion for flexible bodies has the same format as structural vibration.

Using the equivalent static load method, we can derive the equivalent static load using the finite element method. The equation of the motion of a structure under a dynamic load can be written as Eq. (4), where the damping effect is ignored. Rearranging Eq. (4) yields

$$\mathbf{K}_a(\mathbf{r}_a)\mathbf{r}_a = \mathbf{q}(\mathbf{r}_a, t) - \mathbf{M}\ddot{\mathbf{r}}_a \quad (5)$$

or

$$\mathbf{K}_a(\mathbf{r}_a)\mathbf{r}_a = \mathbf{f}_{eq} \quad (6)$$

$$\mathbf{f}_{eq} = \mathbf{q}(\mathbf{r}_a, t) - \mathbf{M}\ddot{\mathbf{r}}_a \quad (7)$$

According to definition of equivalent static load method, Eq. (7) is the equivalent static load at time  $t$ .

Note that the equivalent static load can be described by the external load and inertial force. Thus the equivalent static load is an implicit function of the design variables relating to size, and even though the external force is applied to a single point of a structure, the equivalent static load is applied to all degrees of freedom of the structure.

From Eq. (7), the equivalent static load can only be obtained after performing a transient analysis of the structure to calculate the known displacement fields using an equivalent static load. From the analysis viewpoint, the equivalent static load seems to be useless. However, the goal of the equivalent static load is not to predict displacement induced by a dynamic load but to regenerate the known displacement during an optimization procedure, which will be presented in this work. In other words, the equivalent static load is not an analysis-oriented load but a design-oriented load.

The external loads and inertial forces in the FEM model can be derived from a flexible body dynamic analysis performed by ADAMS software. By using the equivalent static method, we need to build a series of structural equilibrium equations, and then the time-dependent constraint is replaced by the point wise constraints. In such cases, we need to deal with a large number of structural optimization calculations, so we introduce the fast reanalysis method to solve the sequential structural equations with a minimum of computational time. When flexible bodies move at very high speed, the inertial forces will be large and the changes to the structure would alter the performance of the structure significantly, hence the reanalysis and sensitivity analysis methods for traditional structural modification will give diverge or inaccurate results.

### 3. Fast solution method

Given an equivalent structure at position  $p$  with stiffness matrix  $\mathbf{K}_p$  and the equivalent static load vector  $\mathbf{f}_p^{eq}$ , the displacements  $\mathbf{r}_p$  are computed by the equilibrium equation

$$\mathbf{K}_p\mathbf{r}_p = \mathbf{f}_p^{eq} \quad (8)$$

It is assumed that the stiffness matrix  $\mathbf{K}_p$  is given from the initial analysis in the decomposed form

$$\mathbf{K}_p = \mathbf{U}_p^T \mathbf{U}_p \quad (9)$$

where  $\mathbf{U}_p$  is an upper triangular matrix. It is assumed that the degrees of freedom for the next position do not change, so that the finite element equation of the next position  $i$  of the structure becomes

$$\mathbf{K}_{p+i} \mathbf{r}_{p+i} = \mathbf{f}_{p+i}^{eq} \quad (i = 1, 2, \dots) \quad (10)$$

where  $\mathbf{K}_{p+i}$  is the stiffness matrix of the equivalent structure of the next  $i$ th position and vector  $\mathbf{r}_{p+i}$  and  $\mathbf{f}_{p+i}^{eq}$  are the corresponding displacement and load vectors.

The Eq. (10) of the equivalent structure of the next  $i$ th position can be expressed as

$$(\mathbf{K}_p + \Delta \mathbf{K}_{p+i}) \mathbf{r}_{p+i} = \mathbf{f}_p^{eq} + \Delta \mathbf{f}_{p+i}^{eq} \quad (i = 1, 2, \dots) \quad (11)$$

where  $\Delta \mathbf{K}_{p+i}$  is the increment of the stiffness matrix and  $\Delta \mathbf{f}_{p+i}^{eq}$  is the increment of the equivalent static load respectively.

The goal is to find an efficient and accurate approximation of the modified displacements  $\mathbf{r}_{p+i}$  due to the stiffness and equivalent static load changes in the topological modifications without solving Eq. (11). Once the displacements have been evaluated, the explicit stress-displacement relationships can readily determine the stresses

$$\boldsymbol{\sigma} = \mathbf{S} \mathbf{r} \quad (12)$$

where  $\mathbf{S}$  is the stress transformation matrix. The above formulation is suitable for the whole series of equivalent static loads. The difference is only in the  $\Delta \mathbf{K}_{p+i}$  term caused by the different positions due to the motion.

The perturbation method studies the system subjected to small changes in its design parameters. Therefore, if the finite element equilibrium equations of position  $p$  is represented by Eq. (10), the problem becomes that to determine  $\mathbf{r}_{p+i}$  when  $\mathbf{K}_p$  and  $\mathbf{f}_p^{eq}$  are perturbed to the form  $(\mathbf{K}_p + \Delta \mathbf{K}_{p+i})$  and  $\mathbf{f}_p^{eq} + \Delta \mathbf{f}_{p+i}^{eq}$  respectively. So, the static displacement analysis problem of the perturbed system can be written as

$$\mathbf{K} \mathbf{r} = \mathbf{f} \quad (13)$$

where

$$\mathbf{K} = \mathbf{K}_p + \Delta \mathbf{K}_{p+i} \quad (14)$$

$$\mathbf{f} = \mathbf{f}_p^{eq} + \Delta \mathbf{f}_{p+i}^{eq} \quad (15)$$

$$\mathbf{r} = \mathbf{r}_0 + \varepsilon \mathbf{r}_1 + \varepsilon^2 \mathbf{r}_2 + \varepsilon^3 \mathbf{r}_3 + \dots \quad (16)$$

Substituting Eqs. (14), (15) into Eq. (13), we can get

$$\mathbf{r}_0 = \mathbf{K}_p^{-1} \mathbf{f}_p^{eq} \quad (17)$$

$$\mathbf{r}_1 = \mathbf{K}_p^{-1} (\Delta \mathbf{f}_{p+i}^{eq} - \Delta \mathbf{K}_{p+i} \mathbf{r}_0) \quad (18)$$

$$\mathbf{r}_2 = -\mathbf{K}_p^{-1} \Delta \mathbf{K}_{p+i} \mathbf{f}_1 \quad (19)$$

....

$$\mathbf{r}_s = -\mathbf{K}_p^{-1} \Delta \mathbf{K}_{p+i} \mathbf{f}_{s-1} \quad (s = 2, 3, 4, \dots) \quad (20)$$

If we consider an infinite sequence  $\{\mathbf{r}_0, \mathbf{r}_1, \mathbf{r}_2, \dots\}$ , and let the  $\mathbf{s}_n$  be the partial sum of the sequence, then we have a new sequence  $\{\mathbf{s}_0, \mathbf{s}_1, \mathbf{s}_2, \dots\}$

$$\mathbf{s}_n = \sum_{i=0}^n \mathbf{r}_i \quad n = 0, 1, 2, \dots \quad (21)$$

$$\mathbf{s} = \lim_{n \rightarrow \infty} \mathbf{s}_n \quad (22)$$

For the sequence  $\{\mathbf{s}_0, \mathbf{s}_1, \mathbf{s}_2, \dots\}$ , we construct an iterative form

$$\boldsymbol{\varepsilon}_{-1}^{(j)} = \mathbf{0} \quad (23)$$

$$\boldsymbol{\varepsilon}_0^{(j)} = \mathbf{s}_j \quad (24)$$

$$\boldsymbol{\varepsilon}_{k+1}^{(j)} = \boldsymbol{\varepsilon}_{k-1}^{(j+1)} + [\boldsymbol{\varepsilon}_k^{(j+1)} - \boldsymbol{\varepsilon}_k^{(j)}]^{-1}, \quad j = 0, 1, \dots, \quad k = 0, j-1 \quad (25)$$

The iteration formula (23)-(25) is similar to the scalar case (Wu *et al.* 2007), and the inverse of a vector  $\boldsymbol{\varepsilon}^{-1}$  defined as

$$\boldsymbol{\varepsilon}^{-1} = \frac{\boldsymbol{\varepsilon}^*}{(\boldsymbol{\varepsilon}^H \boldsymbol{\varepsilon})} = \frac{\boldsymbol{\varepsilon}^*}{\sum_{i=1}^d \varepsilon_i^2} \quad (26)$$

where the asterisk denotes the complex conjugate and  $H$  the Hermitian conjugate. The vector epsilon-algorithm table can be also constructed by Eqs. (8)-(10).

In order to control the accuracy of the Epsilon Algorithm for a vector series, an error value is given to assure the demand of better kinematic and dynamic performance of a high-speed mechanism is met. The error  $e_R$  is evaluated by

$$e_R = 100 \frac{\|\boldsymbol{\varepsilon}_k^{(0)} - \boldsymbol{\varepsilon}_{k-2}^{(0)}\|}{\|\boldsymbol{\varepsilon}_k^{(0)}\|} < T \quad (27)$$

where  $T$  is the given tolerance. If sufficient accuracy is not achieved, then increasing the number of iterations  $j$  will result in more computational cost. The maximum number of iterations is controlled by the user. If the maximum iterative number is reached and no solution is achieved within the tolerances required then the specific equations cannot be solved by the reanalysis method. In this situation, a full solution is required to solve this equation and thus becomes the new original solution for the next equivalent static load of the structure. Fig. 1 shows how the reanalysis method is performed in the mechanism system.  $p_i^0$  is the location where a full solution must be performed, and the next few structural dynamic equations can be solved using the reanalysis method just mentioned. The interval between equations can be determined by the rate of change of the stiffness of the system. The number of reanalysis points can be set in advance or interactively during the solution.

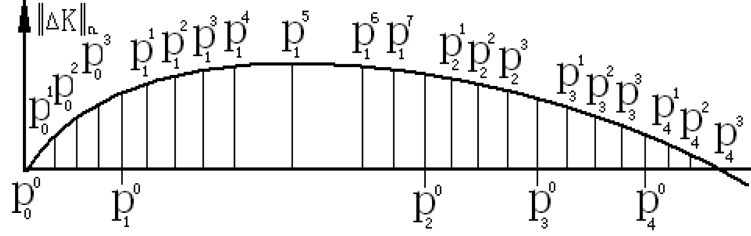


Fig. 1 Indication of solution of serial equivalent static load equations using the reanalysis method

#### 4. Sensitivity analysis

To identify the efficiency of the material, the element sensitivity of the strain energy of elements is introduced. Suppose that an element is removed from the base structure, the element sensitivity of the strain energy is defined by

$$\Delta E_i = E_i^e = \frac{1}{2} \mathbf{r}_i^T \mathbf{K}_i^e \mathbf{r}_i \quad (28)$$

where  $\mathbf{K}_i^e$  is the stiffness matrix of the  $i^{\text{th}}$  element of the structure,  $\mathbf{r}_i$  is the displacement vector of the  $i^{\text{th}}$  element, and  $E_i$  is used to denote the element sensitivity of the strain energy as the result of removing the  $i^{\text{th}}$  element from the reference structure. Using Eq. (28) the all element sensitivity of strain energy of the structures, ( $i = 1, 2, \dots, n$ ), can be obtained. We then define

$$s_i = \frac{1}{c} \sum_{j=1}^c E_{i,j}^e / \max(E_{1,j}^e, E_{2,j}^e, \dots, E_{n,j}^e) \quad (29)$$

where  $n$  is the number of the optimable elements,  $c$  is the number of loadcases,  $\max(E_{1,j}, E_{2,j}, \dots, E_{n,j})$  is the maximum element of the  $j^{\text{th}}$  loadcase;  $s_i$  are the non-dimensional measures for identification of those elements of the structures which need to be kept or removed. Then the criterion for determining the efficiency of material is defined as follows

$$s_i \geq s_k^* \quad (30)$$

where  $s_k^*$  is a threshold value for determining the efficiency of material at  $k^{\text{th}}$  iteration. If the search criterion, Eq. (30), is satisfied, the  $i^{\text{th}}$  element of the structures can be retained; otherwise the  $i^{\text{th}}$  element should be removed.

#### 5. The topological optimization procedure

In the topological optimization of a mechanism's structures, the objective is to find the optimal topology and configuration of material, which minimizes the weight of the structures within global stiffness constraints. The total strain energy is used as the inverse measurement of the global stiffness of the structure. Thus the layout optimization problem of structures can be formulated as follows

$$\begin{aligned} \text{Min } & M \\ \text{s.t. } & E \leq E^* \end{aligned} \quad (31)$$

where  $M$  is the total mass of the structure, and  $E^*$  is a given threshold value,  $E$  is the total strain energy of the structure given by

$$E = \frac{1}{2} \mathbf{r}^T \mathbf{K} \mathbf{r} = \frac{1}{2} \mathbf{f}^T \mathbf{r} \quad (32)$$

where  $\mathbf{K}$  is the stiffness matrix of the structure,  $\mathbf{f}$  and  $\mathbf{r}$  are the load and displacement vectors, respectively.

The procedure of the topological optimization of mechanism's structures is given as follows.

- Step1. Carry out the dynamic analysis of the flexible multibody system, obtain the velocity and acceleration for any position, and export as the load case for each equivalent structural dynamic response.
- Step2. Build finite element models for each position of the motion using the exported boundary condition and formulate a series of static equilibrium equations.
- Step3. Solve the static equilibrium equations using the reanalysis method presented by this paper.
- Step4. Calculate the element sensitivity with regard to the strain energy of each equivalent structure using Eq. (28), and combine all of the element sensitivity results using Eq. (29).
- Step5. Use the search criterion Eq. (30) to decide which elements need to be kept or removed. At the  $k^{\text{th}}$  iteration, the threshold value  $s_k^*$  is determined by

$$s_k^* = s_k^0 + \Delta s_k \quad (33)$$

With this increased rejection threshold value, the elements of the structures can be removed repeatedly until the optimal structures are obtained. The elements are defined as special material instead of being deleted immediately. Therefore, the removed elements can be recovered in later iterations. The procedure requires two parameters to be prescribed. The first is the initial rejection ratio  $s_k^0$  and the second is the increased rate  $\Delta s_k^i$ . Values of  $s_k^0 = 0.01$  and  $\Delta s_k = 0.005$  have been used for many test examples.

- Step6. Update the displacements and element strain energy of structures using the error controlled Epsilon algorithm.

- Step7. Evaluate the total strain energy  $E$  of the structure. Repeat steps 4-6 until the following condition is no longer satisfied, the optimal layout of structures is obtained

$$E \leq E^* \quad (34)$$

or the following condition is satisfied, the the layout optimization of structures is also considered to be obtained.

$$\Delta M < \varepsilon \quad (35)$$

where  $M$  is the mass of removed structures at  $k^{\text{th}}$  iteration, and  $\varepsilon$  is a given small value.

- Step8. Remove the elements by deleting material from the final result, and smooth the boundary using a smoothing method such as the OOSmooth software in Hyperworks.



## 6. Numerical example

During the operation of a die bonder, the bonder (Fig. 2) moves on to the wafer and picks up a die, and then moves back and bonds it onto the leadframe. As the capillary moves at very high speeds, it is necessary to optimally design the mechanism so as to obtain the best dynamic performance. The assembly model of the die bonder with the arm at a  $90^\circ$  drive angle is shown in Fig. 2 and the corresponding assembly finite element model is shown in Fig. 3. The finite element of the arm of the die bonder includes 23153 solid elements and 25744 nodes. The material of structures is aluminum with Young's modulus 71.75 GPa and mass density  $2740 \text{ kg/m}^3$ . During the

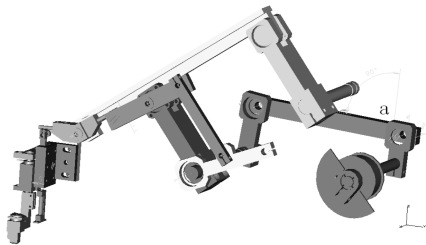


Fig. 2 Assembly of the die bonder

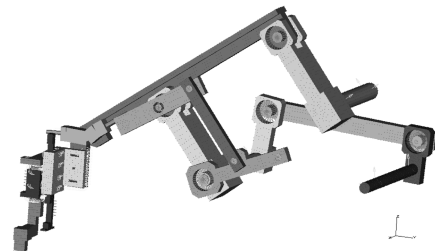


Fig. 3 Finite element Model of the die bonder

Table 1 Frequencies of the first ten modes of the die bonder in different position

| $\alpha$<br>( $^\circ$ ) | Model<br>Freq(Hz) | Mode 2<br>Freq(Hz) | Mode 3<br>Freq(Hz) | Mode 4<br>Freq(Hz) | Mode 5<br>Freq(Hz) | Mode 6<br>Freq(Hz) | Mode 7<br>Freq(Hz) | Mode 8<br>Freq(Hz) | Mode 9<br>Freq(Hz) | Mode 10<br>Freq(Hz) |
|--------------------------|-------------------|--------------------|--------------------|--------------------|--------------------|--------------------|--------------------|--------------------|--------------------|---------------------|
| 0                        | 0.039             | 112.6832           | 117.7509           | 221.7676           | 232.284            | 396.1588           | 510.8086           | 629.2661           | 728.7595           | 884.9595            |
| 10                       | 0.045             | 112.6861           | 185.1197           | 221.2857           | 234.6903           | 396.5793           | 510.5535           | 629.742            | 728.8834           | 884.1193            |
| 20                       | 0.027             | 112.699            | 218.9099           | 227.4708           | 298.0035           | 398.4157           | 509.8322           | 631.1407           | 729.0269           | 880.8168            |
| 30                       | 0.088             | 112.733            | 216.0793           | 228.8021           | 370.9492           | 409.0406           | 509.0369           | 633.3381           | 729.333            | 875.4335            |
| 40                       | 0.175             | 112.8015           | 212.3618           | 229.1991           | 388.5817           | 448.9342           | 508.9621           | 636.1638           | 730.0094           | 868.7857            |
| 50                       | 0.264             | 112.9155           | 208.2248           | 229.6411           | 390.5916           | 481.3221           | 511.8741           | 639.4336           | 731.2723           | 861.8146            |
| 60                       | 0.342             | 113.0796           | 203.9979           | 230.3455           | 390.6403           | 494.6158           | 521.8949           | 642.9725           | 733.3023           | 855.3318            |
| 70                       | 0.412             | 113.2918           | 199.9167           | 231.436            | 390.0111           | 497.7388           | 533.0556           | 646.6244           | 736.2223           | 849.9092            |
| 80                       | 0.473             | 113.5424           | 196.1336           | 232.9899           | 389.0336           | 498.707            | 539.4266           | 650.2569           | 740.0904           | 845.8785            |
| 90                       | 0.528             | 113.8166           | 192.7367           | 235.0379           | 387.8753           | 499.1963           | 540.3386           | 653.7616           | 744.914            | 843.3852            |
| 100                      | 0.584             | 114.0953           | 189.7678           | 237.5547           | 386.6549           | 499.2882           | 535.9728           | 657.0528           | 750.6514           | 842.4299            |
| 110                      | 0.64              | 114.3593           | 187.2372           | 240.4532           | 385.4669           | 498.0139           | 526.9898           | 660.0629           | 757.186            | 842.8739            |
| 120                      | 0.703             | 114.5904           | 185.1348           | 243.5729           | 384.3826           | 490.8803           | 516.9234           | 662.7319           | 764.2741           | 844.4258            |
| 130                      | 0.771             | 114.7728           | 183.4375           | 246.6675           | 383.4332           | 469.3288           | 512.4347           | 665.004            | 771.5254           | 846.682             |
| 140                      | 0.838             | 114.8909           | 182.115            | 249.3502           | 382.4955           | 431.9569           | 511.8312           | 666.8307           | 778.4422           | 849.2025            |
| 150                      | 0.894             | 114.9172           | 181.1338           | 250.7785           | 372.7502           | 385.9856           | 512.2935           | 668.1843           | 784.5007           | 851.5892            |
| 160                      | 0.914             | 114.749            | 180.4611           | 245.8238           | 303.2969           | 382.8814           | 512.9369           | 669.0674           | 789.2372           | 853.5415            |
| 170                      | 0.805             | 113.6145           | 180.0518           | 185.6277           | 266.3805           | 382.5668           | 513.4389           | 669.5137           | 792.3149           | 854.8826            |
| 180                      | 0.039             | 104.3495           | 129.7162           | 179.9479           | 263.5894           | 382.4923           | 513.671            | 669.5726           | 793.5544           | 855.5465            |

motion, the drive angle changes from 0 to 180 degrees. From the modal analysis at different positions, it can be seen that the natural resonance frequencies vary with drive angle (Table 1).

The frequencies of the first ten modes of all positions (with drive angle increments of 10 degrees) are listed in Table 1, and the frequency vs. drive angle for the first ten modes are shown in Fig. 4. These results show that the first mode is rigid motion, and modes 3-7 are more sensitive to drive angle.

In order to reduce the mass of the die bonder such that the inertial force can be reduced, the component structures of the die bonder arm are optimized using the presented methods. The frequencies of the first ten flexible modes during the process of optimization are listed in the Table 2.

In the first iteration, 1334 elements at low sensitivity are modified; note that the rate of change for each component is different. In the second iteration, the material of all the elements are reassigned according to the combined sensitivity, 2050 elements are defined as deleted material, and Young's

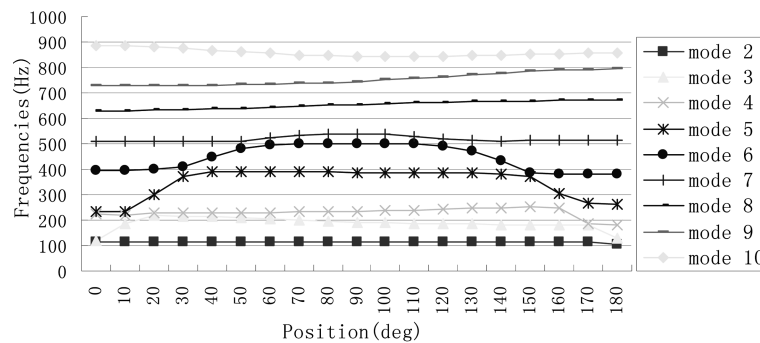


Fig. 4 Frequency vs. drive angle for the first ten modes

Table 2 The frequencies of the first ten modes of the specific position during optimal design

| mode              |                      | Frequencies of structures (Hz) |           |          |          | Change   | Rate of Change |
|-------------------|----------------------|--------------------------------|-----------|----------|----------|----------|----------------|
|                   |                      | Original                       | Step1     | Step2    | Step3    |          |                |
|                   | 2                    | 119.6361                       | 118.2105  | 120.276  | 119.8268 | 0.1907   | 0.16%          |
|                   | 3                    | 204.2693                       | 206.3296  | 202.1958 | 198.1035 | -6.1658  | -3.02%         |
|                   | 4                    | 247.507                        | 252.7171  | 240.7929 | 237.5691 | -9.9379  | -4.02%         |
|                   | 5                    | 328.2244                       | 396.2702  | 323.9933 | 315.9957 | -12.2287 | -3.73%         |
|                   | 6                    | 452.7336                       | 498.1879  | 465.587  | 465.7856 | 13.052   | 2.88%          |
|                   | 7                    | 492.8783                       | 623.4556  | 487.9892 | 483.922  | -8.9563  | -1.82%         |
|                   | 8                    | 623.8224                       | 786.0137  | 610.6964 | 598.4155 | -25.4069 | -4.07%         |
|                   | 9                    | 734.9551                       | 993.0482  | 744.2667 | 742.308  | 7.3529   | 1.00%          |
|                   | 10                   | 928.4488                       | 1099.8681 | 919.6725 | 918.7554 | -9.6934  | -1.04%         |
|                   | 11                   | 1069.6187                      | 1132.44   | 1097.703 | 1106.515 | 36.8963  | 3.45%          |
| Mass<br>(kg)      | Total                | 0.5602                         | 0.5263    | 0.5029   | 0.4876   | -0.0726  | -12.96%        |
|                   | Optimable Structures | 0.3378                         | 0.3035    | 0.2799   | 0.2652   | -0.0726  | -21.49%        |
| Deleting Elements |                      | 0                              | 1334      | 2050     | 2591     | 2591     | 11.19%         |

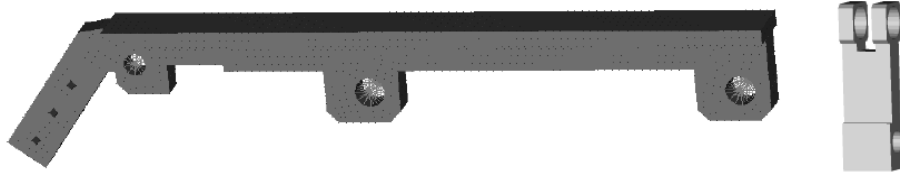


Fig. 5 Original structure of Link 5 and crank 3 of the die bonder

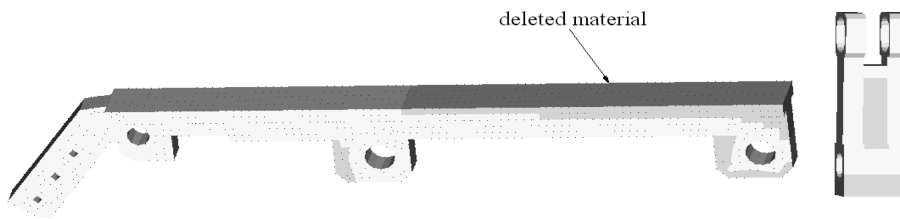


Fig. 6 Removing elements of structures Link 5 and crank 3 of the die bonder

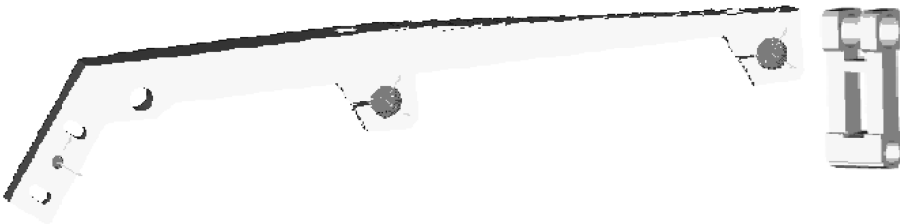


Fig. 7 Optimal structure of Link 5 and crank 3 of the die bonder

Table 3 The comparison of residual vibration and drive torque

|         | Residual vibration          |                 | Drive Torque     |                 |
|---------|-----------------------------|-----------------|------------------|-----------------|
|         | Magnitude ( $\mu\text{m}$ ) | Change Rate (%) | Magnitude (N·mm) | Change Rate (%) |
| Base    | 35.4820                     | -               | 834.9584         | -               |
| Optimal | 28.4600                     | -19.7903        | 672.1309         | -19.5013        |

modulus and mass density are within 1% of the normal material (Fig. 6). In the final optimal design where the elements are completely removed (Fig. 7), the frequencies of the vibration are changed very little, but the mass of the structure is reduced 21.49%. These results show that the presented approach is efficient and easy to carry out.

In order to show the efficiency of optimization, the residual vibration (Fig. 8) and drive torque (Fig. 9) are compared between the base structure and the optimal structure obtained by the presented method. Note that the base structure is the most stiffer structure of the possible design space. It can be seen (shown in Table 3) that the maximum residual vibration magnitude is smaller (decreased 19.7903%) but nearly the same due to the stiffness constraints. Meanwhile, the drive torque became smaller due to the reduced mass thus needing less power (reduced 19.5013%) for the same motion profile of the die bonder.

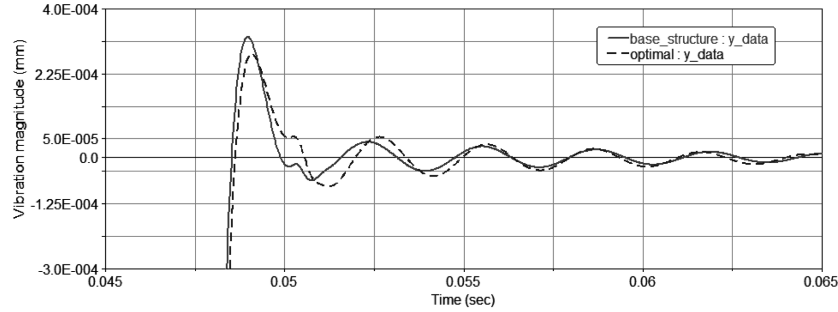


Fig. 8 The residual vibration magnitude curve

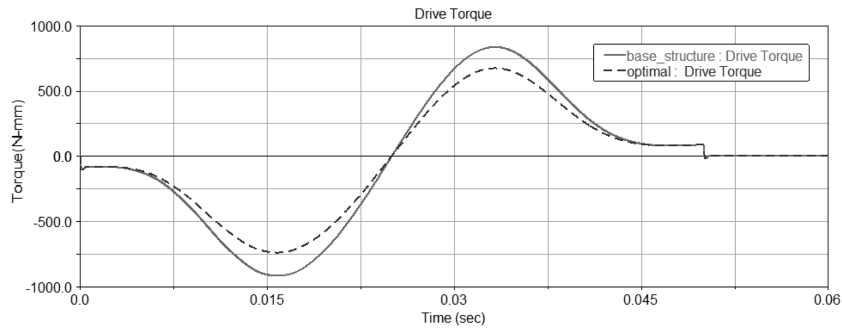


Fig. 9 The driven torque curve

## 7. Conclusions

In this paper, a topological optimization method for flexible multi-body dynamic system has been presented. In this process, the actual boundary conditions of the control system and the approximate stiffness curve of mechanism are calculated from a flexible multi-body dynamics simulation, and the governing equations of the mechanism are divided into separated finite element equilibrium equations for different positions according to the stiffness curve. For efficiency, the static reanalysis method is utilized to solve these equations, a non-dimensional measurement is introduced to identify the efficiency of the elements, and an improved ESO method is used to solve the optimal design problem. The presented method is easy to implement with a general finite element system and convenient to use in various engineering problems. The present method was implemented for the topological optimization of a die bonder. From the example shown in Table 2, it can be seen that only three iterations were needed for obtaining the optimal configuration of material. The flexible multibody simulation shows that the residual vibration can be restrained by stiffness constraints, and the driven power can be minimized due to the lower inertial force using an updated lightweight design.

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