

An original HSDT for free vibration analysis of functionally graded plates

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Abstract. This work presents a free vibration analysis of functionally graded plates by employing an original high order shear deformation theory (HSDT). This theory use only four unknowns, which is even less than the classical HSDT. The equations of motion for the dynamic analysis are determined via the Hamilton's principle. The original kinematic allows obtaining interesting equations of motion. These equations are solved analytically via Navier procedure. The accuracy of the proposed solution is checked by comparing it with other closed form solutions available in the literature.

Keywords: vibration; functionally graded plate; plate theory

1. Introduction

Functionally graded materials (FGMs) are considered as a type of advanced composite material which was originally proposed in Japan (Bever and Duwez 1972, Koizumi 1993). The interesting advantage of this type of composites is the continuity over a desired direction through a structural component (shell, plate or beam) (Kar and Panda 2015, Bennai *et al.* 2015, Belkorissat *et al.* 2015, Ait Atmane *et al.* 2015, Bakora and Tounsi 2015, Barati and Shahverdi 2016, Ahouel *et al.* 2016, Beldjelili *et al.* 2016). Some kinds of conventional composites suffer in continuity within the thickness direction; such discontinuity can be attenuated by a gradual and smooth variation of mechanical characteristics across the thickness of the structural element as in FGMs. Moreover, FGMs allow us to have high thermal and toughness mechanical characteristics, via a mixing for example ceramic and metal.

For the structural applications generally, functionally graded (FG) plates are employed for which the properties change along the thickness. The FG plates find their applications in many engineering industries such as aerospace, nuclear and biomedical. This increase in engineering applications of FGMs has attracted the attention of many scientists. Several research works are found for the structural investigations of FG plates by employing elasticity solution as well as various HSDTs. The exact 3D solution by Kashtalyan (2004), Karami *et al.* (2017) give benchmark for the investigation. Proposing exact solution is a complex process and moreover it is possible only for a few special cases. Hence it is interesting to develop HSDTs which facilitates to propose solution in

much simpler way and with reasonable accuracy. Most of the HSDTs are two dimensional and suppose plane stress condition which neglects transverse normal stress. The most fundamental deformation theory is classical plate theory (CPT) which is valid for thin plates, as it does not consider shear deformation. This may lead to inaccurate results for thick plates, which is thicker than 1/20 of its larger span (Chi and Chung 2006). This model has been implemented by several scientists to investigate the mechanical response of FG plates (Abrate 2008, Ghannadpour and Alinia 2006, Cheng and Batra 2000). The problem encountered in CPT was overcome by first order shear deformation theory (FSDT) developed by Reissner (1945), Mindlin (1951) which takes into consideration the influence of transverse shear deformation. According to FSDT, the transverse strains are considered to be uniform within the thickness of plate which is unrealistic and requires a shear correction factor to compensate transverse shear strains on the top and bottom of the plate (AddaBedia *et al.* 2015, Hadji *et al.* 2016, Boudierba *et al.* 2016). Praveen and Reddy (1998), Chinosi and Croce (2007), Singha *et al.* (2011), Alieldin *et al.* (2011), Wen and Aliabadi (2012), Castellazzi *et al.* (2013), Meksi *et al.* (2015), Bellifa *et al.* (2016) employed FSDT for the bending and vibration analysis of FG plates. Abdelbari *et al.* (2016) presented an efficient and simple shear deformation theory for free vibration of FG rectangular plates on Winkler-Pasternak elastic foundations. Abdelhak *et al.* (2016) studied thermal buckling response of FG sandwich plates with clamped boundary conditions. For avoiding the use of shear correction coefficients, several HSDT, such as, the third-order shear deformation theory (TSDT) (Reddy 2000, Hosseini-Hashemi *et al.* 2011, Qian *et al.* 2004, Akavci 2006, Tounsi *et al.* 2013, Zidi *et al.* 2014, Bourada *et al.* 2016, Merdaci *et al.* 2016, Raminnea *et al.* 2016, Fahsi *et al.* 2017), the sinusoidal shear deformation theory (SSDT) (Neves *et al.* 2012, Boudierba *et al.* 2013, Fekrar *et al.* 2014, Ait Amar Meziiane *et al.* 2014,

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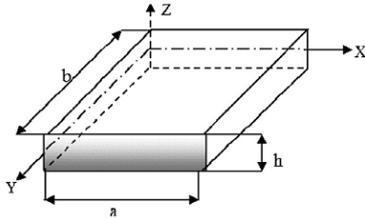


Fig. 1 Geometry of functionally graded plate

Attia *et al.* 2015, AitYahia *et al.* 2015, Hamidi *et al.* 2015, Houari *et al.* 2016, Draïche *et al.* 2016, El-Haina *et al.* 2015, Besseghier *et al.* 2017, Khetir *et al.* 2017), the hyperbolic shear deformation theory (Belabed *et al.* 2014, Hebali *et al.* 2014, Akavci and Tanrikulu 2015, Mahi *et al.* 2015, Saidi *et al.* 2016; Akavci 2016, El-Hassar *et al.* 2016, Mouaïci *et al.* 2016, Bennoun *et al.* 2016, Saidiet *et al.* 2016, Bousahla *et al.* 2014, 2016, Chikh *et al.* 2017, Abualnour *et al.* 2018), the exponential shear deformation theory (Klouche *et al.* 2017) and the zeroth-order shear deformation theory (Bounouara *et al.* 2016, Bellifa *et al.* 2017a) have been developed. It should be noted that many studies are presented in literature to present HSDTs for composite structures as well as graded CNT structures such as (Mehar *et al.* 2017a, b, c, d, Hirwani *et al.* 2017a, b, c; Mehar and Panda 2017, Kar *et al.* 2016, Sahoo *et al.* 2016a, b, Kar and Panda 2016a, b, Mahapatra and Panda 2016, Mehar and Panda 2016, Mehar *et al.* 2016, Mahapatra *et al.* 2016a, b, Katariya and Panda 2016, Sahoo *et al.* 2017, Singh and Panda 2017, Kar and Panda 2015, Singh and Panda 2015, Mahapatra and Panda 2015, Panda and Katariya 2015, Mahapatra *et al.* 2015, Panda and Mahapatra 2014).

In the present paper, the free vibration analysis of FG plates is investigated. The theory contains only four variables, which is less than others HSDTs. The mechanical properties of the plates are supposed to vary in the thickness direction according to a power law distribution in terms of the volume fractions of the constituents. The interesting plate equations of motion for the free vibration analysis are determined through the Hamilton's principle. These equations are then solved using Navier's procedure. The accuracy of the results of this theory is verified by comparing with other HSDTs available in the literature.

2. Analytical modeling

The geometry of a FG plate is as presented in Fig. 1. The dimensions of the plates are $a \times b \times h$, where 'a' is the length, 'b' is width and 'h' is thickness of the plate. The gradation of material characteristics is in the thickness direction with metal and ceramic being the typical constituents. Aluminum/Alumina (Al/Al₂O₃), Aluminum/Zirconia (Al/ZrO₂) and Aluminum/Silicon nitride (Al/Si₃N₄) are the examples of the FG plate.

2.1 Material variation laws

The material properties of FG plate such as the Young's

Table 1 Material properties used in the FG plate

Properties	Metal		Ceramic	
	Aluminum (Al)	Alumina (Al ₂ O ₃)	Zirconia (ZrO ₂)	Silicon nitride (Si ₃ N ₄)
E (GPa)	70	380	200	322.2
ρ (kg/m ³)	2702	3800	5700	2370

modulus E and the mass density ρ are considered to vary continuously within the thickness of the plate according to the power law variation as follows (Bessaim *et al.* 2013, Meradjah *et al.* 2015, Benbakhti *et al.* 2016, Benchohra *et al.* 2017, Benahmed *et al.* 2017, Bouafia *et al.* 2017)

$$E(z) = E_m + (E_c - E_m) \left(\frac{1}{2} + \frac{z}{h} \right)^{p_1}, \quad (1)$$

$$\rho(z) = \rho_m + (\rho_c - \rho_m) \left(\frac{1}{2} + \frac{z}{h} \right)^{p_2}$$

Where (E_c, ρ_c) and (E_m, ρ_m) are the corresponding properties of the ceramic and metal, respectively, and ρ_1, ρ_2 are constants. The Poisson's ratio ν is considered to be constant and equal to 0.3 throughout the analyses (Bourada *et al.* 2012, Taïbi *et al.* 2015, Zemri *et al.* 2015, Bourada *et al.* 2015, Larbi Chaht *et al.* 2015, Laoufi *et al.* 2016, Benferhat *et al.* 2016, Boukhari *et al.* 2016, Chikh *et al.* 2016, Hebali *et al.* 2016, Meksi *et al.* 2017, Mouffoki *et al.* 2017, Zidi *et al.* 2017, Menasria *et al.* 2017). The value of p (p_1 or p_2) equal to zero represents a fully ceramic plate and infinite p , a fully metallic plate. The distribution of the composition of ceramics and metal is linear for $p=1$. Typical values for metal and ceramics used in the FG plate are listed in Table 1.

2.2 Displacement base field

In this work, further simplifying considerations are made to the classical HSDT so that the number of variables is reduced. The displacement field of the classical HSDT is given by (Barka *et al.* 2016, Bellifa *et al.* 2017).

$$u(x, y, z, t) = u_0(x, y, t) - z \frac{\partial w_0}{\partial x} + f(z) \varphi_x(x, y, t) \quad (2a)$$

$$v(x, y, z, t) = v_0(x, y, t) - z \frac{\partial w_0}{\partial y} + f(z) \varphi_y(x, y, t) \quad (2b)$$

$$w(x, y, z, t) = w_0(x, y, t) \quad (2c)$$

Where $u_0; v_0; w_0, \varphi_x, \varphi_y$ are five unknown displacements of the mid-plane of the plate, $f(z)$ denotes shape function representing the variation of the transverse shear strains and stresses within the thickness. By considering that $\varphi_x = \int \theta(x, y) dx$ and $\varphi_y = \int \theta(x, y) dy$, the displacement field of the present theory can be written in a simpler form as

$$u(x, y, z, t) = u_0(x, y, t) - z \frac{\partial w_0}{\partial x} + k_1 f(z) \int \theta(x, y, t) dx \quad (3a)$$

$$v(x, y, z, t) = v_0(x, y, t) - z \frac{\partial w_0}{\partial y} + k_2 f(z) \int \theta(x, y, t) dy \quad (3b)$$

$$w(x, y, z, t) = w_0(x, y, t) \tag{3c}$$

In this work, the present original HSDT is obtained by setting

$$f(z) = \frac{1}{2} h \tanh\left(\frac{2z}{h}\right) - \frac{4}{3} \frac{z^3}{h^2 \cosh(1)^2} \tag{4}$$

It can be seen that the displacement field in Eq. (3) introduces only four unknowns (u_0, v_0, w_0 and θ). The constants k_1 and k_2 depends on the geometry.

In this work, the displacement field considers terms with integrals instead of derivatives, as often employed for example in the CPT or many HSDTs. Such strategy can be important since HSDTs with reduced number of variables can be proposed by utilizing the present foundation.

2.3 Kinematic relations and constitutive relations

In the derivation of the necessary equations, small strains are considered (i.e., displacements and rotations are small, and obey Hooke's law). The linear strain relations derived from the displacement model of Eq. (3) are as follows

$$\begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \end{Bmatrix} = \begin{Bmatrix} \varepsilon_x^0 \\ \varepsilon_y^0 \\ \gamma_{xy}^0 \end{Bmatrix} + z \begin{Bmatrix} k_x^b \\ k_y^b \\ k_{xy}^b \end{Bmatrix} + f(z) \begin{Bmatrix} k_x^s \\ k_y^s \\ k_{xy}^s \end{Bmatrix}, \tag{5}$$

$$\begin{Bmatrix} \gamma_{yz} \\ \gamma_{xz} \end{Bmatrix} = g(z) \begin{Bmatrix} \gamma_{yz}^0 \\ \gamma_{xz}^0 \end{Bmatrix},$$

Where

$$\begin{Bmatrix} \varepsilon_x^0 \\ \varepsilon_y^0 \\ \gamma_{xy}^0 \end{Bmatrix} = \begin{Bmatrix} \frac{\partial u_0}{\partial x} \\ \frac{\partial v_0}{\partial x} \\ \frac{\partial u_0}{\partial y} + \frac{\partial v_0}{\partial x} \end{Bmatrix}, \quad \begin{Bmatrix} k_x^b \\ k_y^b \\ k_{xy}^b \end{Bmatrix} = \begin{Bmatrix} -\frac{\partial^2 w_0}{\partial x^2} \\ -\frac{\partial^2 w_0}{\partial y^2} \\ -2 \frac{\partial^2 w_0}{\partial x \partial y} \end{Bmatrix}, \tag{6a}$$

$$\begin{Bmatrix} k_x^s \\ k_y^s \\ k_{xy}^s \end{Bmatrix} = \begin{Bmatrix} k_1 \theta \\ k_2 \theta \\ k_1 \frac{\partial}{\partial y} \int \theta dx + k_2 \frac{\partial}{\partial x} \int \theta dy \end{Bmatrix},$$

$$\begin{Bmatrix} \gamma_{yz}^0 \\ \gamma_{xz}^0 \end{Bmatrix} = \begin{Bmatrix} k_2 \int \theta dy \\ k_1 \int \theta dx \end{Bmatrix},$$

And

$$g(z) = \frac{df(z)}{dz} \tag{6b}$$

The analytical solution of this theory can be solved by a Navier type solution. The following relations can be obtained

$$\frac{\partial}{\partial y} \int \theta dx = A' \frac{\partial^2 \theta}{\partial x \partial y}, \quad \frac{\partial}{\partial x} \int \theta dy = B' \frac{\partial^2 \theta}{\partial x \partial y},$$

$$\int \theta dx = A' \frac{\partial \theta}{\partial x}, \quad \int \theta dy = B' \frac{\partial \theta}{\partial y} \tag{7}$$

Where the coefficients A' and B' (defined according to the type of solution adopted), k_1 and k_2 are expressed as follows

$$A' = -\frac{1}{\alpha^2}, \quad B' = -\frac{1}{\beta^2}, \quad k_1 = \alpha^2, \quad k_2 = \beta^2 \tag{8}$$

Note that α and β are terms related to the Navier type solution defined in Eq. (23).

For the FG plates, the stress-strain relationships can be written as

$$\begin{Bmatrix} \sigma_x \\ \sigma_y \\ \tau_{xy} \\ \tau_{yz} \\ \tau_{xz} \end{Bmatrix} = \begin{bmatrix} C_{11} & C_{12} & 0 & 0 & 0 \\ C_{12} & C_{22} & 0 & 0 & 0 \\ 0 & 0 & C_{66} & 0 & 0 \\ 0 & 0 & 0 & C_{44} & 0 \\ 0 & 0 & 0 & 0 & C_{55} \end{bmatrix} \begin{Bmatrix} \varepsilon_x \\ \varepsilon_y \\ \gamma_{xy} \\ \gamma_{yz} \\ \gamma_{xz} \end{Bmatrix} \tag{9}$$

where $(\sigma_x, \sigma_y, \tau_{xy}, \tau_{yz}, \tau_{xz})$ and $(\varepsilon_x, \varepsilon_y, \gamma_{xy}, \gamma_{yz}, \gamma_{xz})$ are the stress and strain components, respectively. Using the material properties defined in Eq. (1), stiffness coefficients, C_{ij} , can be given as

$$C_{11} = C_{22} = \frac{E(z)}{1-\nu^2}, \quad C_{12} = \frac{\nu E(z)}{1-\nu^2},$$

$$C_{44} = C_{55} = C_{66} = \frac{E(z)}{2(1+\nu)}, \tag{10}$$

2.3 Equations of motion

Hamilton's principle is employed for the free vibration problem and defined as follows (Hanifi Hachemi Amar *et al.* 2017)

$$0 = \int_0^t (\delta U - \delta K) dt \tag{11}$$

Where δU is the variation of strain energy; and δK is the variation of kinetic energy.

The variation of strain energy of the plate is given by

$$\begin{aligned} \delta U &= \int_V [\sigma_x \delta \varepsilon_x + \sigma_y \delta \varepsilon_y + \tau_{xy} \delta \gamma_{xy} + \tau_{yz} \delta \gamma_{yz} + \tau_{xz} \delta \gamma_{xz}] dV \\ &= \int_A [N_x \delta \varepsilon_x^0 + N_y \delta \varepsilon_y^0 + N_{xy} \delta \gamma_{xy}^0 + M_x^b \delta k_x^b + M_y^b \delta k_y^b + M_{xy}^b \delta k_{xy}^b \\ &\quad + M_x^s \delta k_x^s + M_y^s \delta k_y^s + M_{xy}^s \delta k_{xy}^s + S_{yz}^s \delta \gamma_{yz}^s + S_{xz}^s \delta \gamma_{xz}^s] dA = 0 \end{aligned} \tag{12}$$

Where A is the top surface and the stress resultants $N, M,$ and S are defined by

$$(N_i, M_i^b, M_i^s) = \int_{-h/2}^{h/2} (1, z, f) \sigma_i dz, \quad (i = x, y, xy)$$

$$\text{and } (S_{xz}^s, S_{yz}^s) = \int_{-h/2}^{h/2} g(\tau_{xz}, \tau_{yz}) dz \tag{13}$$

The variation of kinetic energy of the plate can be written as

$$\begin{aligned}
\delta K &= \int_V [\dot{u} \delta \dot{u} + \dot{v} \delta \dot{v} + \dot{w} \delta \dot{w}] \rho(z) dV \\
&= \int_A [I_0 \dot{u}_0 \delta \dot{u}_0 + \dot{v}_0 \delta \dot{v}_0 + \dot{w}_0 \delta \dot{w}_0] \\
&\quad - I_1 \left(\dot{u}_0 \frac{\partial \delta \dot{w}_0}{\partial x} + \frac{\partial \dot{w}_0}{\partial x} \delta \dot{u}_0 + \dot{v}_0 \frac{\partial \delta \dot{w}_0}{\partial y} + \frac{\partial \dot{w}_0}{\partial y} \delta \dot{v}_0 \right) \\
&\quad + J_1 \left((k_1 A') \left(\dot{u}_0 \frac{\partial \delta \dot{\theta}}{\partial x} + \frac{\partial \dot{\theta}}{\partial x} \delta \dot{u}_0 \right) + (k_2 B') \left(\dot{v}_0 \frac{\partial \delta \dot{\theta}}{\partial y} + \frac{\partial \dot{\theta}}{\partial y} \delta \dot{v}_0 \right) \right) \\
&\quad + I_2 \left(\frac{\partial \dot{w}_0}{\partial x} \frac{\partial \delta \dot{w}_0}{\partial x} + \frac{\partial \dot{w}_0}{\partial y} \frac{\partial \delta \dot{w}_0}{\partial y} \right) \\
&\quad + K_2 \left((k_1 A')^2 \left(\frac{\partial \dot{\theta}}{\partial x} \frac{\partial \delta \dot{\theta}}{\partial x} \right) + (k_2 B')^2 \left(\frac{\partial \dot{\theta}}{\partial y} \frac{\partial \delta \dot{\theta}}{\partial y} \right) \right) \\
&\quad - J_2 \left((k_1 A') \left(\frac{\partial \dot{v}_0}{\partial x} \frac{\partial \delta \dot{\theta}}{\partial x} + \frac{\partial \dot{\theta}}{\partial x} \frac{\partial \delta \dot{v}_0}{\partial x} \right) + (k_2 B') \left(\frac{\partial \dot{v}_0}{\partial y} \frac{\partial \delta \dot{\theta}}{\partial y} + \frac{\partial \dot{\theta}}{\partial y} \frac{\partial \delta \dot{v}_0}{\partial y} \right) \right) dA
\end{aligned} \tag{14}$$

Where dot superscript convention indicates the differentiation with respect to the time variable t ; and (I_i, J_i, K_i) are mass inertias expressed by

$$(I_0, I_1, I_2) = \int_{-h/2}^{h/2} (1, z, z^2) \rho(z) dz \tag{15a}$$

$$(J_1, J_2, K_2) = \int_{-h/2}^{h/2} (f, z, f^2) \rho(z) dz \tag{15b}$$

Employing the displacement strain relations (Eqs. (5) and (6)) and stress strain relations (Eq. (9)); applying integrating by parts and the fundamental lemma of variational calculus; and collecting the coefficients of δu_0 , δv_0 , δw_0 , $\delta \theta$ in Eq. (11), the equations of motion are obtained

$$\begin{aligned}
\delta u_0 : \frac{\partial N_x}{\partial x} + \frac{\partial N_{xy}}{\partial y} &= I_0 \ddot{u}_0 - I_1 \frac{\partial \ddot{w}_0}{\partial x} + k_1 A' J_1 \frac{\partial \ddot{\theta}}{\partial x} \\
\delta v_0 : \frac{\partial N_{xy}}{\partial x} + \frac{\partial N_y}{\partial y} &= I_0 \ddot{v}_0 - I_1 \frac{\partial \ddot{w}_0}{\partial y} + k_2 B' J_1 \frac{\partial \ddot{\theta}}{\partial y} \\
\delta w_0 : \frac{\partial^2 M_x^b}{\partial x^2} + 2 \frac{\partial^2 M_{xy}^b}{\partial x \partial y} + \frac{\partial^2 M_y^b}{\partial y^2} &= I_0 \ddot{w}_0 + I_1 \left(\frac{\partial \ddot{u}_0}{\partial x} + \frac{\partial \ddot{v}_0}{\partial y} \right) - I_2 \nabla^2 \ddot{w}_0 \\
&\quad + J_2 \left(k_1 A' \frac{\partial^2 \ddot{\theta}}{\partial x^2} + k_2 B' \frac{\partial^2 \ddot{\theta}}{\partial y^2} \right) \\
\delta \theta : -k_1 M_x^s - k_2 M_y^s - (k_1 A' + k_2 B') \frac{\partial^2 M_{xy}^s}{\partial x \partial y} &+ k_1 A' \frac{\partial \delta S_x^s}{\partial x} + k_2 B' \frac{\partial \delta S_y^s}{\partial y} \\
&= -J_1 \left(k_1 A' \frac{\partial \ddot{u}_0}{\partial x} + k_2 B' \frac{\partial \ddot{v}_0}{\partial y} \right) - K_2 \left((k_1 A')^2 \frac{\partial^2 \ddot{\theta}}{\partial x^2} + (k_2 B')^2 \frac{\partial^2 \ddot{\theta}}{\partial y^2} \right) \\
&\quad + J_2 \left(k_1 A' \frac{\partial^2 \ddot{w}_0}{\partial x^2} + k_2 B' \frac{\partial^2 \ddot{w}_0}{\partial y^2} \right)
\end{aligned} \tag{16}$$

Substituting Eq. (5) into Eq. (9) and the subsequent results into Eqs. (13), the stress resultants are obtained in terms of strains as following compact form

$$\begin{Bmatrix} N \\ M^b \\ M^s \end{Bmatrix} = \begin{bmatrix} A & B & B^s \\ B & D & D^s \\ B^s & D^s & H^s \end{bmatrix} \begin{Bmatrix} \varepsilon \\ k^b \\ k^s \end{Bmatrix}, \quad S = A^s \gamma, \tag{17}$$

In which

$$N = \{N_x, N_y, N_{xy}\}^t, M^b = \{M_x^b, M_y^b, M_{xy}^b\}^t, M^s = \{M_x^s, M_y^s, M_{xy}^s\}^t, \tag{18a}$$

$$\begin{aligned}
\varepsilon &= \{\varepsilon_x^0, \varepsilon_y^0, \gamma_{xy}^0\}^t, \quad k^b = \{k_x^b, k_y^b, k_{xy}^b\}^t, \\
k^s &= \{k_x^s, k_y^s, k_{xy}^s\}^t,
\end{aligned} \tag{18b}$$

$$\begin{aligned}
A &= \begin{bmatrix} A_{11} & A_{12} & 0 \\ A_{12} & A_{22} & 0 \\ 0 & 0 & A_{66} \end{bmatrix}, \quad B = \begin{bmatrix} B_{11} & B_{12} & 0 \\ B_{12} & B_{22} & 0 \\ 0 & 0 & B_{66} \end{bmatrix}, \\
D &= \begin{bmatrix} D_{11} & D_{12} & 0 \\ D_{12} & D_{22} & 0 \\ 0 & 0 & D_{66} \end{bmatrix},
\end{aligned} \tag{18c}$$

$$\begin{aligned}
B^s &= \begin{bmatrix} B_{11}^s & B_{12}^s & 0 \\ B_{12}^s & B_{22}^s & 0 \\ 0 & 0 & B_{66}^s \end{bmatrix}, \quad D^s = \begin{bmatrix} D_{11}^s & D_{12}^s & 0 \\ D_{12}^s & D_{22}^s & 0 \\ 0 & 0 & D_{66}^s \end{bmatrix}, \\
H^s &= \begin{bmatrix} H_{11}^s & H_{12}^s & 0 \\ H_{12}^s & H_{22}^s & 0 \\ 0 & 0 & H_{66}^s \end{bmatrix},
\end{aligned} \tag{18d}$$

$$S = \{S_{xz}^s, S_{yz}^s\}^t, \quad \gamma = \{\gamma_{xz}^0, \gamma_{yz}^0\}^t, \quad A^s = \begin{bmatrix} A_{44}^s & 0 \\ 0 & A_{55}^s \end{bmatrix}, \tag{18e}$$

And stiffness components are expressed as

$$\begin{Bmatrix} A_{11} & B_{11} & D_{11} & B_{11}^s & D_{11}^s & H_{11}^s \\ A_{12} & B_{12} & D_{12} & B_{12}^s & D_{12}^s & H_{12}^s \\ A_{66} & B_{66} & D_{66} & B_{66}^s & D_{66}^s & H_{66}^s \end{Bmatrix} \tag{19a}$$

$$= \int_{-h/2}^{h/2} C_{11} (1, z, z^2, f(z), z f(z), f^2(z)) \begin{Bmatrix} 1 \\ v \\ \frac{1-v}{2} \end{Bmatrix} dz$$

$$(A_{22}, B_{22}, D_{22}, B_{22}^s, D_{22}^s, H_{22}^s) = (A_{11}, B_{11}, D_{11}, B_{11}^s, D_{11}^s, H_{11}^s) \tag{19b}$$

$$A_{44} = A_{55} = \int_{-h/2}^{h/2} C_{44} [g(z)]^2 dz, \tag{19c}$$

By substituting Eq. (17) into Eq. (16), the equations of motion can be expressed in terms of displacements (u_0, v_0, w_0, θ) and the appropriate equations take the form

$$\begin{aligned}
&A_{11} d_{11} u_0 + A_{66} d_{22} u_0 + (A_{12} + A_{66}) d_{12} v_0 - B_{11} d_{111} w_0 \\
&\quad - (B_{12} + 2B_{66}) d_{122} w_0 + (B_{66}^s (k_1 A' + k_2 B')) d_{122} \theta \\
&\quad + (B_{11}^s k_1 + B_{12}^s k_2) d_1 \theta = I_0 \ddot{u}_0 - I_1 d_1 \ddot{w}_0 + J_1 A' k_1 d_1 \ddot{\theta},
\end{aligned} \tag{20a}$$

$$\begin{aligned}
&A_{22} d_{22} v_0 + A_{66} d_{11} v_0 + (A_{12} + A_{66}) d_{12} u_0 - B_{22} d_{222} w_0 \\
&\quad - (B_{12} + 2B_{66}) d_{112} w_0 + (B_{66}^s (k_1 A' + k_2 B')) d_{112} \theta \\
&\quad + (B_{22}^s k_2 + B_{12}^s k_1) d_2 \theta = I_0 \ddot{v}_0 - I_1 d_2 \ddot{w}_0 + J_1 B' k_2 d_2 \ddot{\theta},
\end{aligned} \tag{20b}$$

$$\begin{aligned}
&B_{11} d_{111} u_0 + (B_{12} + 2B_{66}) d_{122} u_0 + (B_{12} + 2B_{66}) d_{112} v_0 \\
&\quad + B_{22} d_{222} v_0 - D_{11} d_{1111} w_0 - 2(D_{12} + 2D_{66}) d_{1122} w_0 \\
&\quad - D_{22} d_{2222} w_0 + (D_{11}^s k_1 + D_{12}^s k_2) d_{11} \theta \\
&\quad + 2(D_{66}^s (k_1 A' + k_2 B')) d_{1122} \theta + (D_{12}^s k_1 + D_{22}^s k_2) d_{22} \theta \\
&= I_0 \ddot{w}_0 + I_1 (d_1 \ddot{u}_0 + d_2 \ddot{v}_0) - I_2 (d_{11} \ddot{w}_0 + d_{22} \ddot{w}_0) \\
&\quad + J_2 (k_1 A' d_{11} \ddot{\theta} + k_2 B' d_{22} \ddot{\theta})
\end{aligned} \tag{20c}$$

Table 2 Comparison of fundamental frequency parameter $\Omega = \omega h \sqrt{\rho_c / E_c}$ simply supported Al/Al₂O₃ square plates

a/h	Theories	Power indices (p)					
		0	0.5	1	4	10	∞
20	Benachour <i>et al.</i> (2011)	0.01480	0.01254	0.01130	0.00980	0.00940	-
	Hosseini-Hashemi <i>et al.</i> (2010)	0.01480	0.01281	0.01150	0.01013	0.00963	-
	Zhao <i>et al.</i> (2009)	0.01464	0.01241	0.01118	0.00970	0.00931	-
	Present	0.01480	0.01254	0.01130	0.00980	0.00940	-
10	Benachour <i>et al.</i> (2011)	0.05769	0.04900	0.04417	0.03804	0.03635	0.02936
	Matsunaga (2008)	0.05777	0.04917	0.04427	0.03811	0.03642	0.02933
	Hosseini-Hashemi <i>et al.</i> (2010)	0.05769	0.04920	0.04454	0.03825	0.03627	0.02936
	Zhao <i>et al.</i> (2009)	0.05673	0.04818	0.04346	0.03757	0.03591	-
	Matsunaga (2008)	0.06382	0.05429	0.04889	0.04230	0.04047	-
	Present	0.05769	0.04900	0.04417	0.03804	0.03635	0.02936
5	Benachour <i>et al.</i> (2011)	0.2112	0.1806	0.1628	0.1375	0.1300	0.1075
	Matsunaga (2008)	0.2121	0.1819	0.1640	0.1383	0.1306	0.1077
	Hosseini-Hashemi <i>et al.</i> (2010)	0.2112	0.1806	0.1650	0.1371	0.1304	0.1075
	Zhao <i>et al.</i> (2009)	0.2055	0.1757	0.1587	0.1356	0.1284	-
	Matsunaga (2008)	0.2334	0.1997	0.1802	0.1543	0.1462	-
	Present	0.2112	0.1806	0.1628	0.1375	0.1300	0.1075

$$\begin{aligned}
 & - (B_{11}^s k_1 + B_{12}^s k_2) d_1 u_0 - (B_{66}^s (k_1 A' + k_2 B')) d_{122} u_0 \\
 & - (B_{66}^s (k_1 A' + k_2 B')) d_{112} v_0 - (B_{12}^s k_1 + B_{22}^s k_2) d_2 v_0 \\
 & + (D_{11}^s k_1 + D_{12}^s k_2) d_{11} w_0 + 2 (D_{66}^s (k_1 A' + k_2 B')) d_{1122} w_0 \\
 & + (D_{12}^s k_1 + D_{22}^s k_2) d_{22} w_0 - H_{11}^s k_1^2 \theta - H_{22}^s k_2^2 \theta \\
 & - 2 H_{12}^s k_1 k_2 \theta - ((k_1 A' + k_2 B')^2 H_{66}^s) d_{1122} \theta \\
 & + A_{44}^s (k_2 B')^2 d_{22} \theta + A_{55}^s (k_1 A')^2 d_{11} \theta \\
 & = -J_1 (k_1 A' d_1 \ddot{u}_0 + k_2 B' d_2 \ddot{v}_0) + J_2 (k_1 A' d_{11} \ddot{w}_0 + k_2 B' d_{22} \ddot{w}_0) \\
 & - K_2 ((k_1 A')^2 d_{11} \ddot{\theta} + (k_2 B')^2 d_{22} \ddot{\theta})
 \end{aligned} \tag{20d}$$

Where d_{ij} , d_{ijl} and d_{ijlm} are the following differential operators

$$\begin{aligned}
 d_{ij} &= \frac{\partial^2}{\partial x_i \partial x_j}, & d_{ijl} &= \frac{\partial^3}{\partial x_i \partial x_j \partial x_l}, \\
 d_{ijlm} &= \frac{\partial^4}{\partial x_i \partial x_j \partial x_l \partial x_m}, & d_i &= \frac{\partial}{\partial x_i}, \quad (i, j, l, m = 1, 2).
 \end{aligned} \tag{21}$$

3. Solution procedure

For the analytical solution of the partial differential equations (Eq. (20)), the Navier's procedure is used. Using this method, the solution of the displacement variables satisfying the boundary conditions can be written in the following Fourier series

$$\begin{Bmatrix} u_0 \\ v_0 \\ w_0 \\ \theta \end{Bmatrix} = \sum_{m=1}^{\infty} \sum_{n=1}^{\infty} \begin{Bmatrix} U_{mn} e^{i\omega t} \cos(\alpha x) \sin(\beta y) \\ V_{mn} e^{i\omega t} \sin(\alpha x) \cos(\beta y) \\ W_{mn} e^{i\omega t} \sin(\alpha x) \sin(\beta y) \\ X_{mn} e^{i\omega t} \sin(\alpha x) \sin(\beta y) \end{Bmatrix} \tag{22}$$

Where ω is the frequency of free vibration of the plate, $\sqrt{-1}$ the imaginary unit.

Where

$$\alpha = m\pi / a \quad \text{and} \quad \beta = n\pi / b \tag{23}$$

Substituting Eq. (22) in Eq. (20), the following equations are obtained

$$\begin{pmatrix} S_{11} & S_{12} & S_{13} & S_{14} \\ S_{12} & S_{22} & S_{23} & S_{24} \\ S_{13} & S_{23} & S_{33} & S_{34} \\ S_{14} & S_{24} & S_{34} & S_{44} \end{pmatrix} - \omega^2 \begin{pmatrix} m_{11} & 0 & m_{13} & m_{14} \\ 0 & m_{22} & m_{23} & m_{24} \\ m_{13} & m_{23} & m_{33} & m_{34} \\ m_{14} & m_{24} & m_{34} & m_{44} \end{pmatrix} \begin{Bmatrix} U_{mn} \\ V_{mn} \\ W_{mn} \\ X_{mn} \end{Bmatrix} = \begin{Bmatrix} 0 \\ 0 \\ 0 \\ 0 \end{Bmatrix} \tag{24}$$

Where

$$\begin{aligned}
 S_{11} &= -(A_{11} \alpha^2 + A_{66} \beta^2), & S_{12} &= -\alpha \beta (A_{12} + A_{66}), \\
 S_{13} &= \alpha (B_{11} \alpha^2 + B_{12} \beta^2 + 2B_{66} \beta^2), \\
 S_{14} &= \alpha (k_1 B_{11}^s + k_2 B_{12}^s - (k_1 A' + k_2 B') B_{66}^s \beta^2), \\
 S_{22} &= -(A_{66} \alpha^2 + A_{22} \beta^2), & S_{23} &= \beta (B_{22} \beta^2 + B_{12} \alpha^2 + 2B_{66} \alpha^2), \\
 S_{24} &= \beta (k_2 B_{22}^s + k_1 B_{12}^s - (k_1 A' + k_2 B') B_{66}^s \alpha^2) \\
 S_{33} &= -(D_{11} \alpha^4 + 2(D_{12} + 2D_{66}) \alpha^2 \beta^2 + D_{22} \beta^4), \\
 S_{34} &= -k_1 (D_{11}^s \alpha^2 + D_{12}^s \beta^2) + 2(k_1 A' + k_2 B') D_{66}^s \alpha^2 \beta^2 \\
 &\quad - k_2 (D_{22}^s \beta^2 + D_{12}^s \alpha^2) \\
 S_{44} &= -k_1 (H_{11}^s k_1 + H_{12}^s k_2) - (k_1 A' + k_2 B')^2 H_{66}^s \alpha^2 \beta^2 \\
 &\quad - k_2 (H_{12}^s k_1 + H_{22}^s k_2) - (k_1 A')^2 A_{55}^s \alpha^2 - (k_2 B')^2 A_{44}^s \beta^2
 \end{aligned}$$

$$\begin{aligned}
 m_{11} &= -I_0, & m_{13} &= \alpha I_1, & m_{14} &= -J_1 k_1 A' \alpha, & m_{22} &= -I_0, \\
 m_{23} &= \beta I_1, & m_{24} &= -k_2 B' \beta J_1, & m_{33} &= -I_0 - I_2 (\alpha^2 + \beta^2) \\
 m_{34} &= J_2 (k_1 A' \alpha^2 + k_2 B' \beta^2), \\
 m_{44} &= -K_2 ((k_1 A')^2 \alpha^2 + (k_2 B')^2 \beta^2)
 \end{aligned} \tag{25}$$

4. Numerical results and discussions

In this section, various numerical examples are presented and discussed for checking the accuracy of the present HSDT in predicting the dynamic behaviors of simply supported FG plates. For the verification purpose,

Table 3 Comparison of fundamental frequency parameter ($\Omega = \omega a^2 \sqrt{\rho_c / E_c} / h$) for simply supported square FG plates when $h/a=0.1$

FGMs Theories	Power indices (p)							
	0	0.5	1	2	5	8	10	
Al/Al ₂ O ₃	Benachour <i>et al.</i> (2011)	5.7694	4.9000	4.4166	4.0057	3.7660	3.6831	3.6357
	Hosseini-Hashemi <i>et al.</i> (2010)	5.7693	4.9207	4.4545	4.0063	3.7837	3.6830	3.6277
	Zhao <i>et al.</i> (2009)	5.6763	4.8209	4.3474	3.9474	3.7218	3.6410	3.5923
	Present	5.7694	4.9000	4.4166	4.0057	3.7660	3.6831	3.6357
Al/ZrO ₂	Benachour <i>et al.</i> (2011)	5.7694	5.4380	5.3113	5.2923	5.3904	5.3950	5.3783
	Matsunaga (2008)	5.7769	-	5.3216	-	-	-	-
	Hosseini-Hashemi <i>et al.</i> (2010)	5.7693	5.3176	5.2532	5.3084	5.2940	5.2312	5.1893
	Zhao <i>et al.</i> (2009)	5.6763	5.1105	4.8713	4.6977	4.5549	4.4741	4.4323
	Present	5.7694	5.4380	5.3113	5.2923	5.3904	5.3950	5.3783

Table 4 Comparison of fundamental frequency parameter $\Omega = \omega h^2 \sqrt{\rho_m / E_m}$ for simply supported square FG plates

Theories	$p=1$			$h/a=0.2$		
	$h/a=0.05$	$h/a=0.1$	$h/a=0.2$	$p=2$	$p=3$	$p=5$
Benachour <i>et al.</i> (2011)	0.0158	0.0618	0.2270	0.2249	0.2255	0.2266
Matsunaga (2008)	0.0158	0.0618	0.2285	0.2264	0.2270	0.2281
Pradyumna and Bandyopadhyay (2008)	0.0157	0.0613	0.2257	0.2237	0.2243	0.2253
Hosseini-Hashemi <i>et al.</i> (2010)	0.0158	0.0611	0.2270	0.2249	0.2254	0.2265
Pradyumna and Bandyopadhyay (2008)	0.0162	0.0633	0.2323	0.2325	0.2334	0.2334
Present	0.0158	0.0618	0.2270	0.2249	0.2255	0.2267

the results obtained by the proposed HSDT are compared with the existing data in the literature and discussed for checking the accuracy of the present HSDT in predicting the dynamic behaviors of simply supported FG plates. For the verification purpose, the results obtained by the proposed HSDT are compared with the existing data in the literature.

Table 2 presented the non-dimensional fundamental frequencies of the simply supported square Al/Al₂O₃ plates for different values of the thickness to length ratios ($h/a=20$; 10 and 5) with $p_1=p_2=p=0, 0.5, 1, 4, 10$, and ∞ . The obtained results are compared with those of Matsunaga (2008), Zhao *et al.* (2009), Hosseini-Hashemi *et al.* (2010), Benachour *et al.* (2011). It should be noted that the results given by Matsunaga (2008) are based on the both FSDT and 2D HSDT; whereas Zhao *et al.* (2009) used the FSDT and employed different values of shear correction factors in their study work. However, the results reported by Hosseini-Hashemi *et al.* (2010) are based on FSDT with a new formula for the shear correction factors. The results given by Benachour *et al.* (2011) are based on refined plate theory (RPT). For convenience in comparison, a novel frequency parameter is defined as $\Omega = \omega h \sqrt{\rho_c / E_c}$. From Table 2, it can be seen that the present results are in excellent agreement with those given by the 2D HSDT (Matsunaga 2008), the RPT (Benachour *et al.* 2011) and the FSDT (Hosseini-Hashemi *et al.* 2010) which use a novel shear correction factors. It is worth noting that all results given on the basis of the FSDT (Matsunaga 2008) are inappropriate since the value of shear correction factor is

considered to be constant ($k^2=1$) for any values of thickness to length ratios and power indices. In addition, the influence of truncated power series to approximate displacement, strain components, and in-plane stress (Matsunaga 2008) on these apparent discrepancies cannot be neglected. The results given by the FSDT (Hosseini-Hashemi *et al.* 2010) are also different from those reported by the proposed theory and by the 2D HSDT (Matsunaga 2008), the RPT (Benachour *et al.* 2011) and the FSDT (Hosseini-Hashemi *et al.* 2010). Another reason of this difference is due to the fact that Zhao *et al.* (2009) utilized a numerical solution (element-free kp -Ritz method) to determine the natural frequencies of the FG plates.

Table 3 shows the comparison of present non-dimensional fundamental frequencies with those obtained with FSDT by Zhao *et al.* (2009) and with higher order theories by Hosseini-Hashemi *et al.* (2010) and Benachour *et al.* (2011) for Al/Al₂O₃ and Al/ZrO₂ square plates square plate with span to thickness ratio 10.

From Table 3, it can be seen that the results of FSDT (Hosseini-Hashemi *et al.* 2010) are in good agreement with both the present theory and the theory presented by Benachour *et al.* (2011) comparatively to those obtained by FSDT Zhao *et al.* (2009). This is due to the fact that the shear correction factor is considered to be constant in the FSDT Zhao *et al.* (2009) ($k^2=5/6$) for any values of power indices, whereas in the FSDT (Hosseini-Hashemi *et al.* 2010) a novel formula for the shear correction factors is supposed taking into account the power indices and the thickness to length ratios. Also, it can be observed that the fundamental frequencies computed by the proposed theory

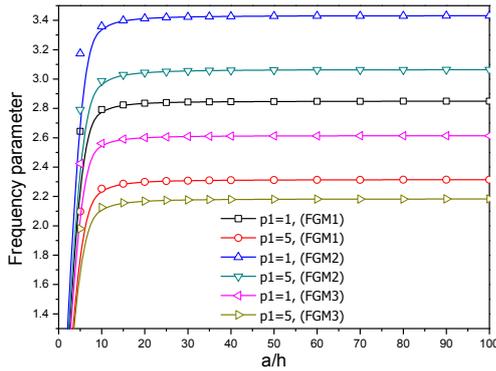


Fig. 2 Variation of frequency parameter with a/h ratio and p_1 index ($a/b=0.5$ and $p_2=1$)

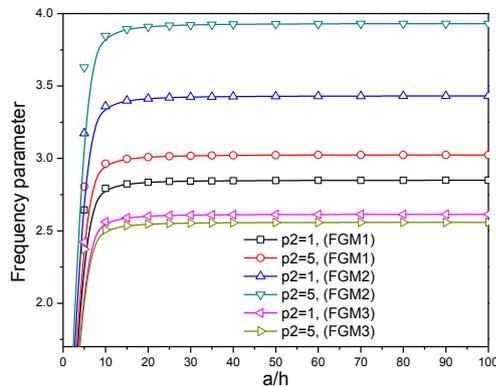


Fig. 3 Variation of frequency parameter with a/h ratio and p_2 index. ($a/b=0.5$ and $p_1=1$)

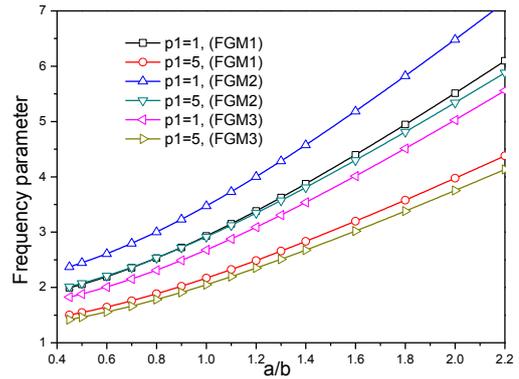


Fig. 4 Variation of frequency parameter with a/b ratio and p_1 index. ($a/h=2$ and $p_2=1$)

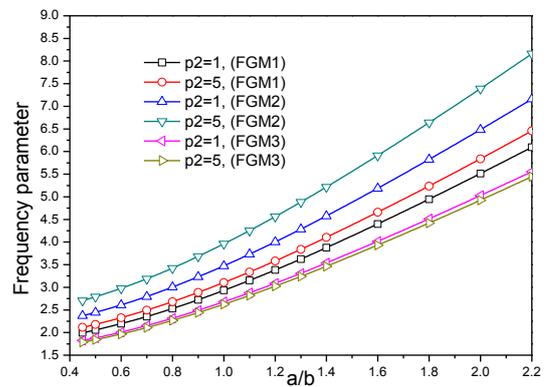


Fig. 5 Variation of frequency parameter with a/h ratio and p_1 index. ($a/h=2$ and $p_1=1$)

demonstrate also a satisfied agreement with 2D HSDT (Matsunaga 2008) for the Al/ZrO₂ square plate.

In Table 4, we present the non-dimensional fundamental frequencies of the Al/ZrO₂ square plate with simply-supported boundary conditions at four edges for $h/a=0.1, 0.2,$ and $1/\sqrt{10}$ when $p_1=p_2=p=0, 1, 2, 3$ and 5 . For convenience in comparison, the non-dimensional fundamental frequency is defined as $\Omega = \omega h^2 \sqrt{\rho_m / E_m}$. Comparing the obtained results with those reported by the 2D HSDT (Matsunaga 2008), FSDT (Hosseini-Hashemi *et al.* 2010) and HSDT of Pradyumna and Bandyopadhyay (2008) and of Benachour *et al.* (2011) demonstrates that all results are in excellent agreement with each other. It is also observed that the proposed analytical method provides the results lower than those reported by the 2D HSDT (Matsunaga 2008) and greater than those given on the basis of the HSDT (Pradyumna and Bandyopadhyay 2008). However, the results computed by the present theory are almost identical to those reported by the FSDT (Hosseini-Hashemi *et al.* 2010) which employ a new shear correction factors. In addition, it should be indicated that the discrepancy between the FSDT (Pradyumna and Bandyopadhyay 2008) and five other theories (i.e., the proposed theory, 2D HSDT (Matsunaga 2008), FSDT (Hosseini-Hashemi *et al.* 2010), and HSDTs of Pradyumna and Bandyopadhyay (2008) and of Benachour *et al.* (2011)) is also considerable.

To make the effects of a/h ratio and power indices more apparent, Figs. 2 and 3 are shown for Aluminum/Alumina (FGM1), Aluminum/Zirconia (FGM2) and Aluminum/ Silicon nitride (FGM3) plates, to show the variation of the non-dimensional fundamental frequency with a/h ratio and p_i ($i=1,2$) power indices, respectively. According to these results the non-dimensional fundamental frequency increases with increasing a/h ratio when $a/h < 20$. The non-dimensional frequency is found to be independent of the length-thickness ratio a/h for $a/h > 20$. It is shown from Fig. 2 that the effect of p_1 is to make the plate stiffer when this gradient index is reduced. However, decreasing the second power index p_2 , makes the plate soften as is presented in Fig. 3. In addition, it is observed that the non-dimensional fundamental frequency is approximately insensitive to p_2 for Aluminum/ Silicon nitride (FGM3) plate.

The dynamic behavior of plate in Ω - a/b plane is presented in Figs. 4 and 5 for FGM 1, FGM 2 and FGM 3. It can be seen that the non-dimensional fundamental frequency increases with increasing the aspect ratio a/b . It is observed from the results that non-dimensional frequency decreases with increasing p_1 however, it increases with increasing p_2 . Again, it can be observed that the non-dimensional frequency is approximately insensitive to p_2 for Aluminum/ Silicon nitride (FGM3) plate.

5. Conclusions

This work presents a free vibration analysis for FG plates by employing an original HSDT with only 4 unknown variables. The nonlinear shear strain function of the theory ensures the accuracy to model the FG plate which depicts non-linear distribution of material characteristics within the thickness. The equations of motion are obtained through the Hamilton's principle. These equations are solved via Navier's procedure. The results were compared with the solutions of several theories. It is concluded that the results of the proposed original HSDT has an excellent agreement with the other theories used for comparison for free vibration problems.

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