# Stability analysis of slopes under groundwater seepage and application of charts for optimization of drainage design 

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#### Abstract

Due to the seepage of groundwater, the resisting force of slopes decreases and the sliding force increases, resulting in significantly reduced slope stability. The instability of most natural slopes is closely related to the influence of groundwater. Therefore, it is important to study slope stability under groundwater seepage conditions. Thus, using a simplified seepage model of groundwater combined with the analysis of stresses on the slip surface, the limit equilibrium (LE) analytical solutions for twoand three-dimensional slope stability under groundwater seepage are deduced in this work. Meanwhile, the general nonlinear Mohr-Coulomb (M-C) strength criterion is adopted to describe the shear failure of a slope. By comparing the results with the traditional LE methods on slope examples, the feasibility of the proposed method is verified. In contrast to traditional LE methods, the proposed method is more suitable for analyzing slope stability under complex conditions. In addition, to facilitate the optimization of drainage design in the slope, stability charts are drawn for slopes with different groundwater tables. Furthermore, the study concluded that: (1) when the hydraulic gradient of groundwater is small, the effect on slope stability is also small for a change in the groundwater table; and (2) compared with a slope without a groundwater table, a slope with a groundwater table has a larger failure range under groundwater seepage.


Keywords: groundwater seepage; two- and three-dimensional slope stability; limit equilibrium; nonlinear MohrCoulomb strength criterion; charts for drainage design

## 1. Introduction

Many factors could affect the slope stability (Zhang et al. 2013, 2015a, b and Peng et al. 2018), and groundwater is one of the most important natural factors (Xu et al. 2010, Piccinini et al. 2014, Luo and Zhang 2016, Zhang et al. 2016). In actual engineering, landslides are often caused as a result of the quick rise and poor drainage of groundwater. With the existence and flow of groundwater, the resisting force of a slope would be not only reduced, but its sliding force is also increased, thereby resulting in a significant decline in slope stability (Vandamme and Zou 2013). With respect to the effect on the resisting force of the slope, the normal stress on the slip surface is reduced under the buoyancy of groundwater so as to decrease the shear capacity of the slip surface. For the sliding force of the slope, the downward trend of the slope becomes more obvious under the infiltration force of groundwater. To study the above effects of groundwater on slope stability, the pore water pressure on the slip surface has been involved to calculate the seepage of groundwater in several methods, such as the limit equilibrium (LE)

[^0]method (Jia et al. 2015, Lu et al. 2015) and limit analysis (LA) method (Zhang et al. 2016).

Using the pore water pressure on the slip surface to show the seepage effect of groundwater is a practical method in the slope stability analysis (Kostic et al. 2015, Jelusic et al. 2016, Li and Yang 2016, Deng et al. 2017). As the calculation of the pore water pressure depends on the shape and location of groundwater table, it is necessary to determine the seepage field of groundwater (Ghiassian and Ghareh 2008). However, the complex groundwater seepage field is usually obtained by numerical simulation with a complicate calculation process, and these obtained results need to be verified by the actual data for judging its applicability. To simplify this analysis process, a previous study used the product of a uniform coefficient (called as the pore water pressure coefficient) and the soil gravity to calculate the pore water pressure on the slip surface (Sun and Zhao 2013). This simplified calculation requires a uniform linear proportional relationship on two vertical heights (one is the vertical height between the groundwater table and slip surface, and another is that between the groundwater table and slope surface), which maybe not accurate with the actual situation. Therefore, it does not accurately reflect the role of groundwater seepage. Here, a simplified model is suggested to establish the groundwater table using some straight lines instead of curved lines so that the analytical solution for the seepage field of groundwater can be easily obtained. In fact, some in-situ monitoring sites are arranged in an actual slope to obtain the groundwater table (Yan et al. 2015, Pirone et al. 2015), and it is approximately determined by linearly connecting these
monitoring sites. Hence, it is feasible to analyze the seepage field of groundwater using this simplified model. Under this simplified model, the calculation of pore water pressure is no longer simply regarded as the product of pore water pressure coefficient and soil gravity, and thus the more reasonable results could be obtained. Moreover, the pore water pressure would be considered as two parts, namely the buoyancy and infiltration force of groundwater, which could be can be calculated analytically in the simplified model. Meanwhile, the location of groundwater table can be described by some parameters in the simplified model of groundwater seepage, such as the highest position of groundwater table.

In addition, the shear strength of soil is also an important factor related to the slope stability. The shear strength of a natural geotechnical body, including the geotechnical body below the groundwater, generally has a nonlinear relationship with the normal stress acting on the geotechnical body (Gao et al. 2016, Wang et al. 2016). It is only an approximate simplification to replace the nonlinear strength criterion with a linear strength criterion, and the problem of slope stability analysis would become easier under the linear strength criterion. As a simple and practical method, the LE method has been widely adopted by designers to analyze slope stability. However, in the traditional LE methods, it is difficult to obtain the analytical solution of slope stability using the nonlinear strength criterion.

In this work, the simplified model of groundwater seepage is adopted to analyze its effect on slope stability with use of the buoyancy and infiltration force of groundwater on the sliding body. Meanwhile, the shear failure of soil on the slip surface is considered to obey the general nonlinear Mohr-Coulomb (M-C) strength criterion. Then, on the basis of the stress analysis on the slip surface, the LE solutions for two-dimensional (2D) and threedimensional (3D) slopes under groundwater seepage are deduced. By comparison and analysis on examples, the feasibility of the proposed method is verified. Furthermore, with the highest position of groundwater table as the variables, the slope stability charts under different groundwater tables are drawn. By applying these charts, the stability of slope under the specify groundwater table could be quickly got, and thereby the drainage design in the slope would be guided to satisfy the requirement of slope safety. Moreover, the optimal parameters for the drainage design could be also obtained from these charts.

## 2. LEM for stability analysis of a slope under groundwater seepage

### 2.1 LEM for 2D slope stability

As shown in Fig. 1, for a slope with a general shape, the groundwater in the slope is assumed to be flow out at the slope toe. For the highest position of groundwater table, it has the horizontal distance $\left(l_{w}\right)$ and vertical distance $\left(h_{w}\right)$ from the slope vertex. Establishing the $x z$ coordinate system with the slope toe as the origin, the equations of the slope surface, slip surface, and groundwater table are $z=g(x), z=$


Fig. 1 Model of stability analysis of a 2D slope under groundwater seepage
$s(x)$, and $z=f(x)$, respectively. $A$ and $B$ are the upper and lower sliding points of slip surface, respectively. In the case of a vertical micro-slice with width $d x$, it is divides into two parts by the groundwater table. For the part above the groundwater table, the forces acting on it include: the gravity $w_{\mathrm{u}} d x$, the horizontal and vertical seismic forces $k_{H} w_{\mathrm{u}} d x$ and $k_{V} w_{\mathrm{u}} d x$, and the horizontal and vertical loads $q_{x} d x$ and $q_{z} d x$ on the slope surface. Compared with the part above the groundwater table, the part below the groundwater table would be affected by the buoyancy and infiltration force of groundwater. Thus, on this part, the forces include: the gravity $w_{\mathrm{b}} d x$, the horizontal and vertical forces $k_{H} w_{\mathrm{b}} d x$ and $k_{V} w_{\mathrm{b}} d x$, the infiltration force $p d x$, and the normal and shear stresses $\sigma d x / \cos \alpha$ and $\tau d x / \cos \alpha$ on slip surface. In the above parameters, $k_{H}$ and $k_{V}$ are the horizontal and vertical seismic force coefficients, respectively, $p$ is the infiltration pressure, $\sigma$ and $\tau$ are the normal and shear stresses on slip surface, respectively, and $\alpha$ is the horizontal inclination angle tangent to the slip surface in the vertical micro-slice.

Here, the slope sliding is considered as the shear failure of slip surface, which is subject to the general nonlinear MC strength criterion. For the traditional LE method, it is difficult to obtain the analytical solution of slope stability under the nonlinear strength criterion. However, the LE stress method, established by Deng et al. (2015, 2016a, 2016b), can be used to analyze the slope stability with the nonlinear strength criterion. Thus, this work derives the LE solution of slope stability under groundwater seepage using the stress assumptions from Deng et al. (2015, 2016a, 2016b).

In the LE stress method, the normal stress $\sigma$ on the slip surface is assumed to be

$$
\begin{equation*}
\sigma=\lambda_{1} \sigma_{0}^{2 D} \tag{1a}
\end{equation*}
$$

$$
\begin{align*}
\sigma_{0}^{2 D}= & \left\{\left[\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{z}+q_{z}\right)\right]- \\
& {\left.\left[k_{H}\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{x}-q_{x}\right] s_{x}\right\} /\left(1+s_{x}^{2}\right) } \tag{1b}
\end{align*}
$$

where $\sigma_{0}{ }^{2 D}$ is the initial normal stress on the 2D slip surface, which is derived according to the force equilibrium conditions with neglecting the increment of inter-slice forces on the two sides of a vertical micro-slice; $\lambda_{1}$, a dimensionless variable, is the correction coefficient of initial normal stress; $w_{\mathrm{u}}=\gamma(g-f) ; w_{\mathrm{b}}=\left(\gamma_{\text {sat }}-\gamma_{\mathrm{w}}\right)(f-s) ; \gamma$ is
the natural unit weight of soil; $\gamma_{\text {sat }}$ is the saturation unit weight of soil; $\gamma_{\mathrm{w}}$ is the unit weight of water; $p_{x}$ and $p_{z}$ are the components of infiltration pressure along the $x$ and $z$
directions, respectively; $\quad p_{x}=\gamma_{w}(f-s) \frac{f_{x}}{\sqrt{1+f_{x}^{2}}} \quad$; $p_{z}=\gamma_{w}(f-s) \frac{f_{x}^{2}}{\sqrt{1+f_{x}^{2}}} ; f_{x}$ is the first derivative of groundwater table equation $f(x)$ with respect to the $x$-axis; $s_{x}$ is the first derivative of slip surface equation $s(x)$ with respect to $x$ axis; and $s_{x}=\tan \alpha$.

In the general nonlinear M-C criterion, the shear strength is

$$
\begin{equation*}
\tau_{f}=c_{0}\left(1+\frac{\sigma}{\sigma_{t}}\right)^{\frac{1}{m}} \tag{2}
\end{equation*}
$$

where $\sigma$ is the normal stress of soil; $\tau_{f}$ is the shear strength of soil under the normal stress $\sigma ; c_{0}$ is the initial cohesion with $c_{0} \geq 0 ; \sigma_{t}$ is the uniaxial tensile strength with $\sigma_{t} \geq 0$; and $m$ is the nonlinear parameter with $m \geq 1$.

For $m=1$ in Eq. (2), it would represent the linear M-C strength criterion. Then, Eq. (2) can be re-expressed by $\tau_{f}=$ $c+\sigma \tan \varphi$, where $c$ is the cohesion of soil, $\varphi$ is the internal friction angle of soil, $c=c_{0}$, and $\tan \varphi=c_{0} / \sigma_{t}$.

When the slope is in the LE state, the slope factor of safety (FOS) is defined as the ratio of the resisting force on slip surface to the sliding force. For a vertical micro-slice, the FOS can be further simplified as the ratio of the shear strength on slip surface to the shear stress. Thus, using Eq. (2), the shear stress on slip surface can be obtained as

$$
\begin{equation*}
\tau=\frac{1}{F_{s}} c_{0}\left(1+\frac{\sigma}{\sigma_{t}}\right)^{\frac{1}{m}} \tag{3}
\end{equation*}
$$

where $\tau$ is the shear stress on slip surface; and $F_{s}$ is the FOS of slope.

From Eq. (3), the first derivative of shear stress $\tau$ with respect to the normal stress $\sigma$ on slip surface can be obtained as

$$
\begin{equation*}
\frac{d \tau}{d \sigma}=\frac{1}{F_{s}} \frac{c_{0}}{m \sigma_{t}}\left(1+\frac{\sigma}{\sigma_{t}}\right)^{\frac{1-m}{m}} \tag{4}
\end{equation*}
$$

Then, expanding Eq. (3) by the Taylor series with the initial normal stress $\sigma_{0}$ as a reference value and substituting Eqs. (4) into the expansion equation, Eq. (3) can be reexpressed as

$$
\begin{align*}
\tau= & \frac{1}{F_{s}} c_{0}\left(1+\frac{\sigma_{0}^{2 D}}{\sigma_{t}}\right)^{\frac{1}{m}}+ \\
& \frac{1}{F_{s}} \frac{c_{0}}{m \sigma_{t}}\left(1+\frac{\sigma_{0}^{2 D}}{\sigma_{t}}\right)^{\frac{1-m}{m}}\left(\sigma-\sigma_{0}^{2 D}\right)+  \tag{5}\\
& H\left(\sigma-\sigma_{0}^{2 D}\right)
\end{align*}
$$

where $H\left(\sigma-\sigma_{0}^{2 D}\right)$ is the higher-order error term.
In Eq. (5), the higher-order error term has correlates with the first two terms on the right-hand side of the equation. Thus, the calculation of $H\left(\sigma-\sigma_{0}{ }^{2 \mathrm{D}}\right)$ can be
replaced by linearly amending this two terms. Then, with substitution of Eq. (1a) and introduction of two new variables, the shear stress is assumed as

$$
\begin{gather*}
\tau=\lambda_{2} \tau_{01}^{2 D}+\lambda_{3} \tau_{02}^{2 D}  \tag{6a}\\
\tau_{01}^{2 D}=c_{0}\left(1+\frac{\sigma_{0}^{2 D}}{\sigma_{t}}\right)^{\frac{1}{m}}  \tag{6b}\\
\tau_{02}^{2 D}=\frac{c_{0}}{m \sigma_{t}}\left(1+\frac{\sigma_{0}^{2 D}}{\sigma_{t}}\right)^{\frac{1-m}{m}} \sigma_{0}^{2 D} \tag{6c}
\end{gather*}
$$

where $\lambda_{2}$ and $\lambda_{3}$ are the correction coefficients of shear stress on the slip surface, both of which are dimensionless variables.

As shown in Fig. 1, the force equilibrium conditions in the $x$ and $z$ directions and the moment equilibrium condition of all the forces about the point $\left(x_{c}, z_{c}\right)$ in the sliding body can be determined as

$$
\begin{align*}
& \int\left[\left(-\sigma s^{\prime}+\tau\right)-k_{H}\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)-p_{x}+q_{x}\right] d x=0  \tag{7a}\\
& \int\left[\left(\sigma+\tau s^{\prime}\right)-\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)-p_{z}-q_{z}\right] d x=0  \tag{7b}\\
& \int\left[\left(-\sigma s^{\prime}+\tau\right)\left(y_{c}-s\right)+\left(\sigma+\tau s^{\prime}\right)\left(x-x_{c}\right)\right] d x- \\
& \int\left[\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{z}+q_{z}\right]\left(x-x_{c}\right) d x- \\
& \int k_{H} w_{\mathrm{u}}\left[z_{c}-\frac{1}{2}(f+g)\right] d x-  \tag{7c}\\
& \int\left(p_{x}+k_{H} w_{\mathrm{b}}\right)\left[z_{c}-\frac{1}{2}(s+f)\right] d x+ \\
& \int q_{x}\left(z_{c}-g\right) d x=0
\end{align*}
$$

By substituting Eqs. (7a)-(7b) into Eq. (7c) and simplifying, the following formula can be obtained as

$$
\begin{align*}
& \int\left[\sigma\left(s s^{\prime}+x\right)+\tau\left(x s^{\prime}-s\right)\right] d x- \\
& \int\left[\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{z}+q_{z}\right] x d x+ \\
& \int\left[\frac{1}{2} k_{H} w_{\mathrm{u}}(f+g)+\frac{1}{2}\left(p_{x}+k_{H} w_{\mathrm{b}}\right)(s+f)\right] d x-  \tag{8}\\
& \int q_{x} g d x=0
\end{align*}
$$

From Eq. (8), it is noted that the choice of the moment centre point in the sliding body has no effect on the establishment of the LE equations.

Then, by substituting Eqs. (1a) and (6a) into Eqs. (7a)(7b) and (8), respectively, the following linear equations for the variables $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ can be obtained as

$$
\begin{gather*}
\sum_{j=1}^{3} a_{i j} \lambda_{j}=b_{i} \quad(i=1,2,3)  \tag{9a}\\
a_{11}=-\int \sigma_{0}^{2 D} s^{\prime} d x \tag{9b}
\end{gather*}
$$

$$
\begin{gather*}
a_{12}=\int \tau_{01}^{2 D} d x  \tag{9c}\\
a_{13}=\int \tau_{02}^{2 D} d x  \tag{9d}\\
a_{21}=\int \sigma_{0}^{2 D} d x  \tag{9e}\\
a_{22}=\int \tau_{01}^{2 D} s^{\prime} d x  \tag{9f}\\
a_{23}=\int \tau_{02}^{2 D} s^{\prime} d x  \tag{9~g}\\
a_{31}=\int \sigma_{0}^{2 D}\left(s s^{\prime}+x\right) d x  \tag{9h}\\
a_{32}=\int \tau_{01}^{2 D}\left(x s^{\prime}-s\right) d x  \tag{9i}\\
a_{33}=\int \tau_{02}^{2 D}\left(x s^{\prime}-s\right) d x  \tag{9j}\\
b_{1}=\int\left[k_{H}\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{x}-q_{x}\right] d x  \tag{9k}\\
b_{2}=\int\left[\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{z}+q_{z}\right] d x  \tag{91}\\
b_{3}=\int\left[\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{z}+q_{z}\right] x d x- \\
\int \frac{1}{2} k_{H} w_{\mathrm{u}}(f+g) d x-  \tag{9m}\\
\int\left[\frac{1}{2}\left(p_{x}+k_{H} w_{\mathrm{b}}\right)(s+f)-q_{x} g\right] d x \\
\end{gather*}
$$

According to Eq. (9), the variables ( $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ ) would be solved. Thereafter, with substitution of these variables into Eqs. (1)-(6), the normal stress $\sigma$ and shear stress $\tau$ on the slip surface can be obtained.

Based on the above definition, the slope FOS is regarded as the ratio of the resisting force on slip surface to the sliding force for the shear failure of a slope. Thus, the FOS of a 2D slope can be solved using the ratio of the total resisting force along the entire slip surface to the total sliding force, and its formula is

$$
\begin{equation*}
F_{s}=\frac{\int\left[c_{0}\left(1+\frac{\sigma}{\sigma_{t}}\right)^{\frac{1}{m}} / \cos \alpha\right] d x}{\int(\tau / \cos \alpha) d x} \tag{10}
\end{equation*}
$$

### 2.2 LEM for 3D slope stability

As shown in Fig. 2(a), in the stability analysis of a 3D slope with slope height $H$ and slope angle $\beta$, the symmetry surface of 3D symmetric sliding body (or the main slip surface of 3D asymmetric sliding body) is called the neutral


Fig. 2 Model of stability analysis of a 3D slope under groundwater seepage
plane. Taking the slope toe in the neutral plane as the origin, an $x y z$ coordinate system is established. Then, $z=g(x, y), z$ $=s(x, y)$, and $z=f(x, y)$ are the equations of slope surface, 3D slip surface, and groundwater table, respectively, where the positive $x$-axis points into the slope, the positive $z$-axis is opposite to gravity, and the positive $y$-axis is determined according to the right hand rule.

In Fig. 2(a) and 2(b), to describe the characteristics of a 3D sliding body, the width of 3D sliding body along the $y$ axis is $W$, and the sliding body is divided into two parts with the neutral plane as the interface. The widths of left and right parts are $w_{1}$ and $w_{2}$, respectively. Additionally, points $A, B, C$ and $D$ are the upper, lower, left, and right border endpoints of 3D sliding body, respectively, where points $A$ and $B$ are in the neutral plane, and the length of
line $C D$ represents the width of 3D sliding body.
In contrast to the simple 2D sliding body, the 3D sliding body slides down substantially along the main sliding direction. For a symmetric sliding body, the main sliding direction is parallel to the symmetry plane (or the neutral plane) because of the symmetry of sliding body. For an asymmetric sliding body, the main sliding direction is not parallel to the neutral plane but is at an angle $\rho$ to the neutral plane. Another coordinate system, i.e., $x^{\prime} y^{\prime} z$, is then established with the main sliding direction of 3D sliding body as the $x^{\prime}$-axis. Meanwhile, the $x^{\prime} y^{\prime} z$ coordinate system has the same origin as the $x y z$ coordinate system, and the positive direction of $y^{\prime}$-axis is also determined by the right hand rule. Moreover, the main sliding direction of 3D sliding body is determined by the angle $\rho$, which can be defined as the inclination angle of the total shear forces on 3D slip surface in the $x z$ plane.

Compared to the original $x y z$ coordinate system, rotating the original coordinate system to align with the main sliding direction of sliding body has the advantage that in the $x^{\prime} y^{\prime} z$ coordinate system, only three global LE conditions of 3D sliding body, i.e., the force equilibrium condition in the $x^{\prime}$ and $z$ directions and the moment equilibrium condition around one point in the $y^{\prime}$ direction, are required to solve the LE stability of 3D slope. Thus, the complexity of 3D slope stability analysis is reduced, and the calculation speed is improved. Furthermore, the stability of 2D and 3D slopes would be solved by the universal LE equations.

As shown in Fig. 2(c), the vertical 3D micro-column $A_{1} B_{1} C_{1} D_{1} A_{2} B_{2} C_{2} D_{2}$ is selected from the 3D sliding body. The micro-column has widths $d x$ and $d y$ in the $x$ and $y$ directions, respectively. $\alpha_{x}$ and $\alpha_{y}$ are the inclination angles of slip surface $A_{1} B_{1} C_{1} D_{1}$ in the $x z$ and $y z$ planes, respectively, where $\alpha_{y}$ is positive in the clockwise direction along the $y$-axis and is otherwise negative. Using these parameters, the area of slip surface $A_{1} B_{1} C_{1} D_{1}$ can be calculated as $\Delta d x d y$, where $\Delta=\sqrt{1+s_{x}^{2}+s_{y}^{2}}, s_{x}=\tan \alpha_{x}$, and $s_{y}=\tan \alpha_{y}$. Consistent with the vertical 2D micro-slice, the vertical 3D micro-column $A_{1} B_{1} C_{1} D_{1} A_{2} B_{2} C_{2} D_{2}$ is also divided into two parts by the groundwater table.

In the vertical micro-column under general conditions, for the part above the groundwater table, the forces acting on it have: the gravity $w_{\mathrm{u}} d x d y$, the horizontal and vertical seismic forces $k_{H} w_{\mathrm{u}} d x d y$ and $k_{V} w_{\mathrm{u}} d x d y$, and the loads $q_{x} d x d y, q_{y} d x d y$, and $q_{z} d x d y$ on the slope surface. Moreover, on the part below the groundwater table, the forces include: the gravity $w_{\mathrm{b}} d x d y$, the horizontal and vertical seismic forces $k_{H} w_{\mathrm{b}} d x d y$ and $k_{V} w_{\mathrm{b}} d x d y$, the infiltration forces $p^{x z} d x d y$ and $p^{y z} d x d y$, the normal stress $\sigma \Delta d x d y$ on slip surface, and the shear stresses $\tau^{\gamma z} \Delta d x d y$ and $\tau^{y z} \Delta d x d y$ on slip surface. In these parameters, $p^{x z}$ and $p^{y z}$ are the components of infiltration pressure in the $x z$ and $y z$ planes. Meanwhile, and $\tau^{x z}$ and $\tau^{y z}$ are the components of shear stress on slip surface in the $x z$ and $y z$ planes, respectively.

Consistent with the 2D slope stability analysis, the LE stress method assumes the normal stress on the 3D slip surface to be

$$
\begin{equation*}
\sigma=\lambda_{1} \sigma_{0}^{3 D} \tag{11a}
\end{equation*}
$$

$$
\begin{align*}
\sigma_{0}^{3 D}= & {\left[\frac{\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{z}^{x z}+p_{z}^{y z}+q_{z}}{\Delta}-\right.} \\
& \frac{k_{H}\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{x}^{x z}-q_{x}}{\Delta} \tan \alpha_{x}+  \tag{11b}\\
& \left.\frac{p_{y}^{y z}+q_{y}}{\Delta} \operatorname{sgn}\left(\alpha_{y}\right) \tan \alpha_{y}\right] / \\
& {\left[n_{\sigma}^{z}-n_{\sigma}^{x} \tan \alpha_{x}-n_{\sigma}^{y} \operatorname{sgn}\left(\alpha_{y}\right) \tan \alpha_{y}\right] }
\end{align*}
$$

where $\sigma_{0}{ }^{3 D}$ is the initial normal stress on 3D slip surface; $n_{\sigma}{ }^{x}, n_{\sigma}{ }^{y}$, and $n_{\sigma}{ }^{z}$ are the $x$-, $y$-, and $z$-axis direction cosines of normal stress on slip surface, respectively; $n_{\sigma}{ }^{x}=-s_{x} / \Delta ; n_{\sigma}{ }^{x y}=-$ $s_{y} / \Delta ; \mathrm{n}_{\sigma}{ }^{z}=1 / \Delta ; p_{x}^{x z}$ and $p_{z}{ }^{x z}$ are the components of infiltration force $p^{x z}$ along the $x$ - and $z$-axis directions, respectively;

$$
p_{x}^{x z}=\gamma_{w}(f-s) \frac{f_{x}}{\sqrt{1+f_{x}^{2}} ;} \quad p_{z}^{x z}=\gamma_{w}(f-s) \frac{f_{x}^{2}}{\sqrt{1+f_{x}^{2}}}
$$

$f_{x}$ is the first partial derivative of groundwater table equation $f(x, y)$ with respect to the $x$-axis; $p_{y}^{y z}$ and $p_{z}^{y z}$ are the components of infiltration force $p^{y z}$ along the $y$ - and $z$ axis directions, respectively;

$$
p_{y}^{y z}=\gamma_{w}(f-s) \frac{f_{y}}{\sqrt{1+f_{y}^{2}} ;} \quad p_{z}^{y z}=\gamma_{w}(f-s) \frac{f_{y}^{2}}{\sqrt{1+f_{y}^{2}}}
$$

and $f_{y}$ is the first partial derivative of groundwater table equation $f(x, y)$ with respect to the $y$-axis.

Similar to the 2D slope, the initial normal stress $\sigma_{0}{ }^{3 D}$ on slip surface, shown in Eq. (11), is derived according to the force equilibrium conditions when the increment of intercolumn forces on the two sides of vertical micro-column is assumed to be zero.

According to the above definition of the angle $\rho$, it can be solved using the shear stresses $\tau^{x z}$ and $\tau^{y z}$ on slip surface. For shear stresses $\tau^{x z}$ and $\tau^{y z}$, they can be also approximately deduced according to the force equilibrium conditions as the increments of inter-column forces on the two sides of vertical micro-column are neglected. However, in the derivation process of shear stresses, the normal stress on the slip surface should be $\sigma$ of Eq. (11a) to make the obtained result more reasonable. Thus, the angle $\rho$ is given as

$$
\begin{align*}
\rho= & \arcsin \left\{-\iint\left(\frac{p_{y}^{y z}+q_{y}}{\Delta}+\lambda_{1} n_{\sigma}^{y} \sigma_{0}^{3 D}\right) \Delta d x d y /\right. \\
& \iint \sqrt{\left[\frac{k_{H}\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+p_{x}^{x z}-q_{x}}{\Delta}-\lambda_{1} n_{\sigma}^{x} \sigma_{0}^{3 D}\right]^{2}+}  \tag{12}\\
& \sqrt{\left.\left(\frac{p_{y}^{y z}+q_{y}}{\Delta}+\lambda_{1} n_{\sigma}^{y} \sigma_{0}^{3 D}\right)^{2} \Delta d x d y\right\}}
\end{align*}
$$

The main sliding direction of a 3D sliding body, on the basis of which the $x^{\prime} y^{\prime} z$ coordinate system is established, is determined by the angle $\rho$. Hence, the angle between the planes $x^{\prime} z$ and $x z$ is also $\rho$. Thereby, $n_{\sigma}{ }^{x^{\prime}}$ and $n_{\sigma}{ }^{y^{\prime}}$, which are the $x^{\prime}$ - and $y^{\prime}$-axis direction cosines of normal stress $\sigma$ on slip surface, respectively, can be calculated as

$$
\begin{gather*}
n_{\sigma}^{x^{\prime}}=n_{\sigma}^{x} \cos \rho+n_{\sigma}^{y} \sin \rho  \tag{13a}\\
n_{\sigma}^{y^{\prime}}=-n_{\sigma}^{x} \sin \rho+n_{\sigma}^{y} \cos \rho \tag{13b}
\end{gather*}
$$

To easily solve for the normal stress $\sigma$ and the slope FOS $\left(F_{s}\right)$, the visual shear force on slip surface is oriented parallel to the main sliding direction of 3D sliding body, i.e., the visual shear stress $\tau$ on slip surface being parallel to the $x^{\prime} z$ plane. This assumption is widely used in the traditional 3D LE methods. Thereby, the $y^{\prime}$-axis direction cosines ( $n_{\tau}^{y^{\prime}}$ ) of the visual shear stress $\tau$ on slip surface can be calculated as

$$
\begin{equation*}
n_{\tau}^{y^{\prime}}=0 \tag{14}
\end{equation*}
$$

For the $x^{\prime}$ - and $z$-axis direction cosines of the visual shear stress $\tau$ on slip surface in a 3D vertical micro-column, they are named by $n_{\tau}^{x^{\prime}}$ and $n_{\tau}^{z}$, respectively. Then, these cosines of the normal stress and visual shear stress ( $n_{\tau}^{x^{\prime}}, n_{\tau}^{y^{\prime}}$, $n_{\tau}^{z}, n_{\sigma}{ }^{x^{\prime}}, n_{\sigma}{ }^{y^{\prime}}$, and $n_{\sigma}{ }^{z}$ ) have the following relationships as

$$
\begin{align*}
& \left(n_{\tau}^{x^{\prime}}\right)^{2}+\left(n_{\tau}^{y^{\prime}}\right)^{2}+\left(n_{\tau}^{z}\right)^{2}=1  \tag{15a}\\
& n_{\sigma}^{x^{\prime}} n_{\tau}^{x^{\prime}}+n_{\sigma}^{y^{\prime}} n_{\tau}^{y^{\prime}}+n_{\sigma}^{z} n_{\tau}^{z}=0 \tag{15b}
\end{align*}
$$

By combining Eqs. (13), (14) and (15), the cosines $n_{\tau}^{x^{\prime}}$ and $n_{\tau}^{z}$ can be respectively solved as

$$
\begin{gather*}
n_{\tau}^{x^{\prime}}=\frac{n_{\sigma}^{z} \sqrt{\left[\left(n_{\sigma}^{x^{\prime}}\right)^{2}+\left(n_{\sigma}^{z}\right)^{2}\right]}}{\left(n_{\sigma}^{x^{\prime}}\right)^{2}+\left(n_{\sigma}^{z}\right)^{2}}  \tag{16a}\\
n_{\tau}^{z}=-\frac{n_{\sigma}^{x^{\prime}} n_{\tau}^{x^{\prime}}}{n_{\sigma}^{z}} \tag{16b}
\end{gather*}
$$

Consistent with the 2 D slope stability analysis, the virtual shear stress $\tau$ in the 3D slope can be also calculated using Eq. (3) for the slope sliding due to shear failure of slip surface. Thereby, the calculation of the virtual shear stress $\tau$ is assumed as

$$
\begin{gather*}
\tau=\lambda_{2} \tau_{01}^{3 D}+\lambda_{3} \tau_{02}^{3 D}  \tag{17a}\\
\tau_{01}^{3 D}=c_{0}\left(1+\frac{\sigma_{0}^{3 D}}{\sigma_{t}}\right)^{\frac{1}{m}}  \tag{17b}\\
\tau_{02}^{3 D}=\frac{c_{0}}{m \sigma_{t}}\left(1+\frac{\sigma_{0}^{3 D}}{\sigma_{t}}\right)^{\frac{1-m}{m}} \sigma_{0}^{3 D} \tag{17c}
\end{gather*}
$$

Thereafter, to facilitate the establishment of the LE equations of a 3D sliding body in the $x^{\prime} y^{\prime} z$ coordinate system, the $x$ and $y$ coordinates in the 3D vertical microcolumn are transformed into their $x^{\prime}$ and $y^{\prime}$ coordinates, which are expressed as

$$
\begin{equation*}
x^{\prime}=x \cos \rho+y \sin \rho \tag{18a}
\end{equation*}
$$

$$
\begin{equation*}
y^{\prime}=-x \sin \rho+y \cos \rho \tag{18b}
\end{equation*}
$$

According to the above analysis, only three LE conditions of a 3D sliding body need to be satisfied in the $x^{\prime} y^{\prime} z$ coordinate system, i.e., the force equilibrium conditions of all forces acting on a 3D sliding body in the $x^{\prime}$ and $z$ directions and their moment equilibrium conditions around the point $\left(x_{c}, y_{c}, z_{c}\right)$ in the $y^{\prime}$ direction. Then, they are given as

$$
\begin{align*}
& \iint\left[\left(\sigma n_{\sigma}^{x^{\prime}}+\tau n_{\tau}^{x^{\prime}}\right) \Delta-k_{H}\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right) \cos \rho\right] d x d y- \\
& \iint\left[\left(p_{x}^{x z}-q_{x}\right) \cos \rho-\left(p_{y}^{v z}+q_{y}\right) \sin \rho\right] d x d y=0  \tag{19a}\\
& \iint\left[\left(\sigma_{\sigma}^{z}+\tau n_{\tau}^{z}\right) \Delta-\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)\right] d x d y- \\
& \iint\left(p_{z}^{x z}+p_{z}^{y z}+q_{z}\right) d x d y=0  \tag{19b}\\
& \iint\left(\sigma_{\sigma}^{z}+\tau m_{\tau}^{z}\right) \Delta\left(x^{\prime}-x_{c}\right) d x d y+ \\
& \iint\left(\sigma_{\sigma}^{x^{\prime}}+m_{\tau}^{x^{\prime}}\right) \Delta\left(z_{c}-s\right) d x d y- \\
& \iint\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)\left(x^{\prime}-x_{c}\right) d x d y- \\
& \iint\left(p_{z}^{x z}+p_{z}^{y z}+q_{z}\right)\left(x^{\prime}-x_{c}\right) d x d y-  \tag{19c}\\
& \iint k_{H} w_{\mathrm{u}} \cos \rho\left[z_{c}-(g+f) / 2\right] d x d y- \\
& \iint\left[\left(k_{H} w_{\mathrm{b}}+p_{x}^{x z}\right) \cos \rho\left[z_{c}-(f+s) / 2\right] d x d y+\right. \\
& \iint p_{y}^{v z} \sin \rho\left[z_{c}-(f+s) / 2\right] d x d y- \\
& \iint\left(q_{x} \cos \rho+q_{y} \sin \rho\right)\left(g-z_{c}\right) d x d y=0
\end{align*}
$$

By substituting Eqs. (19a) and (19b) into Eq. (19c), Eq. (19c) can be simplified into

$$
\begin{align*}
& \iint\left[\sigma\left(n_{\sigma}^{z} x^{\prime}-n_{\sigma}^{x^{\prime}} s\right) \Delta+\tau\left(n_{\tau}^{z} x^{\prime}-n_{\tau}^{x^{\prime}} s\right) \Delta d x d y-\right. \\
& \iint\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right) x^{\prime} d x d y+ \\
& \iint \frac{1}{2} k_{H} w_{\mathrm{u}}(g+f) \cos \rho d x d y+ \\
& \iint \frac{1}{2} k_{H} w_{\mathrm{b}}(f+s) \cos \rho d x d y+  \tag{20}\\
& \iint \frac{1}{2}\left(p_{x}^{x z} \cos \rho-p_{y}^{y z} \sin \rho\right)(f+s) d x d y- \\
& \iint\left(p_{z}^{x z}+p_{z}^{y z}\right) x^{\prime} d x d y- \\
& \iint\left[\left(q_{x} \cos \rho+q_{y} \sin \rho\right) g+q_{z} x^{\prime}\right] d x d y=0
\end{align*}
$$

Consist with the 2D analysis, it is also noted from Eq. (20) that the choice of the moment centre point in the 3D sliding body has no effect on the establishment of the LE equations.

By substituting Eqs. (11a) and (17a) into Eqs. (19a)(19b) and (20), respectively, the linear equations of the three variables $\lambda_{1}-\lambda_{3}$ can be obtained as

$$
\begin{align*}
& \sum_{j=1}^{3} c_{i j} \lambda_{j}=d_{i} \quad(i=1,2,3)  \tag{21a}\\
& c_{11}=\iint \sigma_{0}^{3 D} n_{\sigma}^{x^{\prime}} \Delta d x d y  \tag{21b}\\
& c_{12}=\iint \tau_{01}^{3 D} n_{\tau}^{x^{\prime}} \Delta d x d y  \tag{21c}\\
& c_{13}=\iint \tau_{02}^{3 D} n_{\tau}^{x^{\prime}} \Delta d x d y  \tag{21d}\\
& c_{21}=\iint \sigma_{0}^{3 D} n_{\sigma}^{z} \Delta d x d y  \tag{21e}\\
& c_{22}=\iint \tau_{01}^{3 D} n_{\tau}^{z} \Delta d x d y  \tag{21f}\\
& c_{23}=\iint \tau_{02}^{3 D} n_{\tau}^{z} \Delta d x d y  \tag{21~g}\\
& c_{31}=\iint \sigma_{0}^{3 D}\left(n_{\sigma}^{z} x^{\prime}-n_{\sigma}^{x^{\prime}} s\right) \Delta d x d y  \tag{21h}\\
& c_{32}=\iint \tau_{01}^{3 D}\left(n_{\tau}^{z} x^{\prime}-n_{\tau}^{x^{\prime}} s\right) \Delta d x d y  \tag{21i}\\
& c_{33}=\iint \tau_{02}^{3 D}\left(n_{\tau}^{z} x^{\prime}-n_{\tau}^{x^{\prime}} s\right) \Delta d x d y  \tag{21j}\\
& d_{1}=\iint\left[k_{H}\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right)+\left(p_{x}^{x z}-q_{x}\right)\right] \cos \rho d x d y- \\
& \left.\iint\left(p_{y}^{y z}+q_{y}\right) \sin \rho\right] d x d y  \tag{21k}\\
& d_{2}=\iint\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right) d x d y+ \\
& \iint\left(p_{z}^{x z}+p_{z}^{y z}+q_{z}\right) d x d y  \tag{211}\\
& d_{3}=\iint\left[\left(1-k_{V}\right)\left(w_{\mathrm{u}}+w_{\mathrm{b}}\right) x^{\prime} d x d y-\right. \\
& \iint \frac{1}{2} k_{H} w_{\mathrm{u}}(g+f) \cos \rho d x d y- \\
& \iint \frac{1}{2} k_{H} w_{\mathrm{b}}(f+s) \cos \rho d x d y-  \tag{21m}\\
& \iint \frac{1}{2}\left(p_{x}^{x z} \cos \rho-p_{y}^{y z} \sin \rho\right)(f+s) d x d y+ \\
& \iint\left(p_{z}^{x z}+p_{z}^{y z}\right) x^{\prime} d x d y+ \\
& \iint\left[\left(q_{x} \cos \rho+q_{y} \sin \rho\right) g+q_{z} x^{\prime}\right] d x d y
\end{align*}
$$

By solving Eq. (21), the variables ( $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ ) could be obtained. However, according to Eq. (21), the calculation of the parameters ( $c_{i j}$ and $d_{i}$ ) in the Eq. (21) is related to the angle $\rho$. Therefore, the angle $\rho$ should be determined before
solving for the variables ( $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ ). Thus, let the initial value of angle $\rho$ be $\rho^{(0)}=0$, and these variables ( $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ ) then is solved using Eq. (21). Thereafter, by substituting the obtained $\lambda_{1}$ into Eq. (12), a new $\rho$ would be calculated. If $\left|\rho-\rho^{(0)}\right| \leqq \varepsilon$ (here $\varepsilon=0.01^{\circ}$ ), the obtained values of $\rho$, $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ are the final results. If $\left|\rho-\rho^{(0)}\right|>\varepsilon$, let $\rho^{(0)}=$ $\rho$, and $\rho, \lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ are recalculated. Thereby, the looping iteration of $\rho$ is required to solve for the variables ( $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ ) in the 3D slope stability.

After the final results of $\lambda_{1}, \lambda_{2}$, and $\lambda_{3}$ are obtained, the normal stress $\sigma$ and visual shear stress $\tau$ on 3D slip surface can be calculated by substituting them into Eqs. (11) and (17).

Since the virtual shear stress is assumed to be parallel to the main sliding direction of 3D sliding body, the 3D slope FOS can be also solved using the ratio of the total resisting force on slip surface to the total sliding force, which is consistent with the definition of 2D slope FOS. Then, the 3D slope FOS is calculated as

$$
\begin{equation*}
F_{s}=\frac{\iint c_{0}\left(1+\frac{\sigma}{\sigma_{t}}\right)^{\frac{1}{m}} \Delta d x d y}{\iint \tau \Delta d x d y} \tag{22}
\end{equation*}
$$

## 3. Comparison and analysis of slope examples

### 3.12D slope examples

Taking a 2D homogeneous slope as an example, the influence of groundwater seepage on slope stability is considered. To verify the feasibility of the proposed method, the results of the proposed method are compared with those of the traditional LE methods. The slope height and slope angle are assumed to be $H$ and $\beta$, respectively. Moreover, the natural unit weight of soil, saturation unit weight of soil, and unit weight of water are $\gamma, \gamma_{s a t}$, and $\gamma_{w}$, respectively, where $\gamma_{w}=10 \mathrm{kN} / \mathrm{m}^{3}$. In the slope, the soil parameter are given as $c_{0}, \sigma_{t}$, and $m$ under the nonlinear MC strength criterion, and they would be usually named by the cohesion $c$ and internal friction angle $\varphi$ if the linear M-C strength criterion (i.e., $m=1$ in Eq. (1)) is adopted. Here, for the groundwater table, it is simplified as a straight line, which describes that the groundwater flows from the highest position of groundwater table to the slope toe. $l_{w}$ and $h_{w}$ are the horizontal and vertical distances of the highest position of groundwater table from the slope vertex, respectively. Then, according to different combinations of slope parameters, parameters describing groundwater table, and soil parameters, 32 slope examples are formed. Thus, the stability of these slope examples is analyzed by the traditional LE Swedish method (Fellenius 1936), the Morgenstern-Price (M-P) method (Morgenstern and Price 1965), and the proposed method. The calculated results are listed in Table 1. In Table 1, two kinds of slip surface (i.e., circular slip surface and arbitrary curved slip surface) are adopted, and the generation of the arbitrary curved slip surface is from the method of Deng et al. (2017). It is well

Table 1 Comparison of the results of 2D slope stability analysis

|  | Slopeparameters |  | Groundwater table parameters |  | Soil parameters |  |  |  | Minimum FOS of slope |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Examples |  |  |  |  |  |  | $c_{0}$ | $\arctan \left(c_{0} / \sigma_{t}\right)$ |  | Cir | slip s |  | Arbitrar slip s | y curved urface |
|  | $H$ (m) | $\beta$ | $h_{w}(\mathrm{~m})$ | $l_{w}(\mathrm{~m})$ | ( $\mathrm{kN} / \mathrm{m}^{3}$ ) | $\left(\mathrm{kN} / \mathrm{m}^{3}\right)$ | $\begin{aligned} & (\text { or } c) \\ & (\mathrm{kPa}) \end{aligned}$ | $\left(\right.$ or $\varphi$ ) ( ${ }^{\circ}$ ) |  | Swedish method | $\begin{gathered} \text { M-P } \\ \text { method } \end{gathered}$ | Current method | $\begin{gathered} \mathrm{M}-\mathrm{P} \\ \text { method } \end{gathered}$ | Current method |
| 1 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 20 | 1.0 | 1.309 | 1.544 | 1.558 | 1.535 | 1.496 |
| 2 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 25 | 1.0 | 1.474 | 1.733 | 1.747 | 1.726 | 1.682 |
| 3 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 30 | 1.0 | 1.648 | 1.933 | 1.947 | 1.928 | 1.878 |
| 4 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 35 | 1.0 | 1.838 | 2.151 | 2.163 | 2.147 | 2.091 |
| 5 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 20 | 1.0 | 1.591 | 1.883 | 1.902 | 1.867 | 1.813 |
| 6 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 25 | 1.0 | 1.762 | 2.08 | 2.100 | 2.067 | 2.012 |
| 7 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 30 | 1.0 | 1.942 | 2.288 | 2.307 | 2.277 | 2.219 |
| 8 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 35 | 1.0 | 2.138 | 2.512 | 2.531 | 2.503 | 2.441 |
| 9 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 20 | 1.0 | 1.145 | 1.231 | 1.292 | 1.219 | 1.259 |
| 10 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 25 | 1.0 | 1.258 | 1.362 | 1.425 | 1.351 | 1.391 |
| 11 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 30 | 1.0 | 1.380 | 1.500 | 1.565 | 1.489 | 1.531 |
| 12 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 35 | 1.0 | 1.512 | 1.649 | 1.716 | 1.639 | 1.682 |
| 13 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 20 | 1.0 | 1.433 | 1.532 | 1.614 | 1.511 | 1.563 |
| 14 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 25 | 1.0 | 1.552 | 1.665 | 1.753 | 1.650 | 1.706 |
| 15 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 30 | 1.0 | 1.678 | 1.811 | 1.898 | 1.795 | 1.851 |
| 16 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 35 | 1.0 | 1.813 | 1.966 | 2.055 | 1.951 | 2.009 |
| 17 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 20 | 1.5 | 1.043 | - | 1.254 | - | 1.182 |
| 18 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 25 | 1.5 | 1.130 | - | 1.359 | - | 1.291 |
| 19 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 30 | 1.5 | 1.220 | - | 1.467 | - | 1.395 |
| 20 | 10 | 30 | 2 | 20 | 17.8 | 20 | 25 | 35 | 1.5 | 1.314 | - | 1.579 | - | 1.506 |
| 21 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 20 | 1.5 | 1.323 | - | 1.594 | - | 1.489 |
| 22 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 25 | 1.5 | 1.420 | - | 1.709 | - | 1.607 |
| 23 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 30 | 1.5 | 1.519 | - | 1.826 | - | 1.728 |
| 24 | 10 | 30 | 2 | 20 | 17.8 | 20 | 35 | 35 | 1.5 | 1.621 | - | 1.950 | - | 1.852 |
| 25 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 20 | 1.5 | 0.973 | - | 1.092 | - | 1.044 |
| 26 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 25 | 1.5 | 1.037 | - | 1.169 | - | 1.123 |
| 27 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 30 | 1.5 | 1.103 | - | 1.248 | - | 1.204 |
| 28 | 10 | 45 | 2 | 12 | 17.8 | 20 | 25 | 35 | 1.5 | 1.173 | - | 1.330 | - | 1.285 |
| 29 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 20 | 1.5 | 1.264 | - | 1.413 | - | 1.340 |
| 30 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 25 | 1.5 | 1.334 | - | 1.495 | - | 1.426 |
| 31 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 30 | 1.5 | 1.405 | - | 1.580 | - | 1.515 |
| 32 | 10 | 45 | 2 | 12 | 17.8 | 20 | 35 | 35 | 1.5 | 1.480 | - | 1.670 | - | 1.607 |

known that the Swedish method is a non-rigorous method for being only suitable for a circular slip surface, and the MP method has a rigorous solution without the limitation on the type of slip surface.

Table 1 shows that: (1) compared to the simple LE Swedish method, the proposed method satisfies all the global LE conditions of a sliding body so that a rigorous LE solution is obtained, and it could be also applied with the arbitrary curved slip surface (including the circular slip surface); and (2) the proposed method produces similar results to those of the rigorous LE M-P method, but the proposed method can analyze the slope stability under the
nonlinear strength criterion.

### 3.2 3D slope examples

3D slope example 1: as shown in Fig. 3, is a 3D homogenous slope, with slope height $H=40 \mathrm{~m}$, slope angle $\beta=45^{\circ}$. The soil parameters $\gamma=22 \mathrm{kN} / \mathrm{m}^{3}, \gamma_{\text {sat }}=24 \mathrm{kN} / \mathrm{m}^{3}$, $c=30 \mathrm{kPa}$, and $\varphi=30^{\circ}$ for the slope subject to the linear M-C strength criterion. Moreover, $\gamma_{w}=10 \mathrm{kN} / \mathrm{m}^{3}$ for the unit weight of water. For this slope, a 3D spheroid is assumed as the sliding body. Then, by establishing an $x y z$ coordinate system consistent with Fig. 2, the equation of 3D


Fig. 3 3D slope example 1
Table 2 Comparison of the results of 3D slope stability analysis

| W/H | Slope FOS |  |  |  |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | 3D ordinary column method | 3D simplified Janbu method | 3D Spencer method | 3D M-P method | Improved 3D FOS method | Current method |
| 1 | 1.152 | 1.252 | 1.343 | 1.368 | 1.480 | 1.219 |
| 2 | 1.142 | 1.172 | 1.237 | 1.261 | 1.340 | 1.192 |
| 4 | 1.142 | 1.162 | 1.219 | 1.232 | 1.292 | 1.177 |
| 6 | 1.153 | 1.142 | 1.210 | 1.218 | 1.260 | 1.173 |
| 8 | 1.154 | 1.142 | 1.209 | 1.214 | 1.260 | 1.172 |



Fig. 4 Results of 3D slope stability analysis for a dynamic change in groundwater
ellipsoid is given as

$$
\begin{equation*}
\frac{\left(x-x_{0}\right)^{2}}{a^{2}}+\frac{\left(z-z_{0}\right)^{2}}{a^{2}}+\frac{y^{2}}{b^{2}}=1 \tag{23}
\end{equation*}
$$

where $a$ and $b$ are the shape parameters of 3D ellipsoid, and ( $x_{0}, 0, z_{0}$ ) is the center coordinates of the circular slip surface with radius $a$ in the neutral plane of 3D sliding body.

In Eq. (23), the parameter $b$ can be determined on basic of the parameters $x_{0}, z_{0}, a, H, \beta$, and $W$ (i.e., the width W of 3D sliding body) as

For $x_{0}<-z_{0} \tan \beta$ and $x_{0}<0$,

$$
\begin{equation*}
b=\frac{W}{2 \sqrt{1-\frac{z_{0}^{2}}{a^{2}}}} \tag{24a}
\end{equation*}
$$

For $x_{0}>H /(\sin \beta \cos \beta)-z_{0} \tan \beta$ and $x_{0}>H / \tan \beta$,

$$
\begin{equation*}
b=\frac{W}{2 \sqrt{1-\frac{\left(H-z_{0}\right)^{2}}{a^{2}}}} \tag{24b}
\end{equation*}
$$

For other cases,

$$
\begin{equation*}
b=\frac{W}{2 \sqrt{1-\left(\frac{x_{0}}{a} \tan \beta-\frac{z_{0}}{a}\right)^{2} /\left[1+\tan ^{2} \beta\right]}} \tag{24c}
\end{equation*}
$$

To compare the results with other methods, the slope without the groundwater table is given, and the 3D ellipsoid is also used with the parameters of the sliding body as $x_{0}=$ $40 \mathrm{~m}, z_{0}=40 \mathrm{~m}$, and $a=40 \mathrm{~m}$ at the specific position. Then, the FOS of a 3D slope for $W / H=1,2,4,6$, and 8 is calculated, and the results are listed in Table 2. Table 2 shows that the proposed method has the close results with the 3D simplified Janbu method (Hungr 1989), 3D Spencer method (Zhang 1988), and 3D M-P method (Cheng and Yip 2007), thereby verifying the feasibility of the proposed feasibility. Moreover, similar to the 3D ordinary column method (Hovland 1997), only three LE conditions of a 3D sliding body are adopted in the proposed method. However, the proposed method obtains the global LE conditions of a 3D sliding body from the special Cartesian coordinate system, which is established on the basis of the main sliding


Fig. 5 Contrast on the results of critical slip surface and normal stress on slip surface in 3D slope example 2
direction of 3D sliding body. Thus, the rigorous 3D LE solutions can be solved by the proposed method. In addition, the improved 3D FOS method (Wang and Deng 2003) has a larger FOS than the proposed method for the reason that it considers the appearance of sliding between the adjacent 3D vertical columns. In fact, a relative displacement between the adjacent 3D vertical columns may be not obvious when the slope slides due to shear failure of slip surface. Therefore, compared with the improved 3D FOS method, the results obtained by the proposed method are more conservative.

For the slope with the consideration of groundwater table, the effect of a dynamic change in the groundwater table on slope stability is studied. To easily describe the groundwater seepage in a 3D slope, it is assumed that the groundwater flows out from the slope toe, and the groundwater table between its highest position and the slope toe is simplified into a plane, which is the same with the simplified seepage model for a 2D slope. For the highest position of groundwater table, $l_{w}$ and $h_{w}$ are its horizontal and vertical distances from the slope vertex, respectively. Meanwhile, no hydraulic gradient acts on this simple 3D slope along the $y$-axis direction, that is, the seepage does not occur along the width direction of a 3D sliding body. Then, the above 3 D rotating ellipsoid with the same parameters is also adopted as the sliding body to analyze the slope stability. Thus, when $W=20 \mathrm{~m}$, the curves of the minimum FOS of 3D slope vs. the parameter $\left(l_{w} \tan \beta\right) / H$ are plotted for $h_{w} / H=0.1,0.2,0.3,0.4$, and 0.5 in Fig. 4.

Fig. 4 shows that: (1) the slope stability would be reduced under the groundwater seepage, and the slope tends to be more likely instability with the high groundwater table; and (2) a change in the groundwater table would has the small effect on the slope stability under the condition of the low hydraulic gradient of groundwater, and the hydraulic gradient is expressed by $\left(H-h_{w}\right) /\left(H / \tan \beta+l_{w}\right)$.

3D slope example 2: as shown in Fig. 5, the 3D slope with a weak interlayer has slope height $H=12.25 \mathrm{~m}$ and slope ratio $1: 2$, i.e., slope angle $\beta=26.565^{\circ}$. The weak interlayer is located below the slope toe with the vertical distance $h_{d}=0.75 \mathrm{~m}$, and its thickness is $h=0.5 \mathrm{~m}$. The soil parameters are $\gamma=18.84 \mathrm{kN} / \mathrm{m}^{3}, \gamma_{\text {sat }}=21 \mathrm{kN} / \mathrm{m}^{3}, c=28.5$ kPa , and $\varphi=20^{\circ}$ for the slope subject to the linear M-C strength criterion. $\gamma_{w}=10 \mathrm{kN} / \mathrm{m}^{3}$ for the unit weight of
water. For the weak interlayer, which is also subject to the linear M-C criterion, it has $\gamma=18.84 \mathrm{kN} / \mathrm{m}^{3}, \gamma_{\text {sat }}=21$ $\mathrm{kN} / \mathrm{m}^{3}, c=0.0 \mathrm{kPa}$, and $\varphi=10^{\circ}$. Similar to the 3D slope in example 1 , the 3 D rotating ellipsoid as the sliding body is also adopted in this example. However, the presence of a weak interlayer makes it easier for the slope to slide along its surface. Thus, when the 3D rotating ellipse is below the weak interlayer, it would be replaced by the surface of weak interlayer, thereby forming a 3D composite slip surface. Taking the width of 3D sliding body $W=40 \mathrm{~m}$ as an example shown in Fig. 5, the minimum slope FOS, the critical slip surface, and the normal stress on the slip surface are respectively calculated for the two cases, i.e., the slope without groundwater and the slope under groundwater seepage. For the slope under groundwater seepage, the groundwater table is simplified using the same model as in 3D slope example 1. Here, $h_{w}=0.4 H$ and $l_{w}=1.2 H / \tan \beta$. By analyzing the slope stability, the minimum slope FOS is equal to 2.400 for a slope without a groundwater table, and it is 2.055 for the slope under groundwater seepage conditions. Thus, the slope stability under groundwater seepage is significantly reduced, and it also shows that the groundwater is an important natural factor resulting in slope instability.

Fig. 5 illustrates that the normal stress on the slip surface is obviously reduced by the buoyancy force of groundwater, and as a result the resisting force of slope is also decreased. Meanwhile, the infiltration force of groundwater increases the sliding force of slope so as to reduce significantly the slope stability. In addition, the slope under groundwater seepage has a larger failure range than that without the groundwater table.

## 4. Charts for optimization of the design of slope drainage

The above analysis shows that the slope stability is significantly reduced with the existence of groundwater seepage. Hence, in the actual engineering scenarios under groundwater table, the drainage is required by arranging some pipes (i.e., drainage pipe) in the slope to improve the slope stability. In other words, it is very important for drainage design in slope under the groundwater seepage. For the drainage pipe, the plastic pipe commonly used in


Fig. 6 Stability charts for optimizing slope drainage design with slope angle $\beta=30^{\circ}$
the engineering could be adopted, and some drain holes would be set around its shell. Moreover, the mesh geotextile is wrapped at the end of the pipe and around the shell of the pipe to prevent the silt from entering the pipe. Meanwhile, the drain pipe is arranged obliquely downward within the slope so that the water inside the pipe can quickly flow out under the action of its gravity. With regard to the spacing of the drainage pipes, it can be ensured that the groundwater above the drainage pipes and between the two adjacent drainage pipes should be effectively and promptly discharged.

In addition, compared with the non-linear strength criterion, the linear M-C strength criterion is still widely used by designers for the slope stability analysis. The reason is that only two shear strength parameters included in the linear M-C strength criterion can be quickly and simply obtained by the laboratory experiments and the failure behavior of slope sliding can be also approximately simulated with the use of them. Therefore, it is practical to draw the stability charts of slope with the linear M-C strength criterion for optimizing the design of drainage.

Then, the stability charts of slope are drawn with slope height $H=10$ and slope angle $\beta$. In the slope, the soil parameters are $\gamma=17.8 \mathrm{kN} / \mathrm{m}^{3}$ and $\gamma_{\text {sat }}=20 \mathrm{kN} / \mathrm{m}^{3} \cdot \gamma_{w}=10$ $\mathrm{kN} / \mathrm{m}^{3}$ for the unit weight of water. Here, it is thought that the groundwater table would descend to the designed
position under the work of the drainage system after the drainage pipes are arranged in the slope. Thus, to facilitate designing the position of drainage pipes, the groundwater around the pipes would be drained by them, and finally a new stable groundwater table is formed in the slope. Meanwhile, the new groundwater table can be approximately described using the above simplified model, in which the groundwater undergoes the linear flow from the highest position of groundwater table to the slope toe after drainage. Then, the designed highest position of groundwater table is determined by the parameters $l_{w}$ and $h_{w}$. Thereby, the drainage design of slope is considered as the design of parameters describing the groundwater table, and the purpose of arranging the drainage pipe in the slope is to descend the groundwater table to the designed position.

In the work of Sun and Zhao (2013), the curve of $\tan \varphi$ vs. $c /(\gamma H)$ for the slope in the critical state was plot to form the simple and practical stability charts, and it could be also used to get the stability charts under the different groundwater table for optimizing the drainage design. In the stability charts, a slope in the critical state has the minimum FOS of 1.000 , which is used to judge whether the slope is in a stable state. For example, for one point above the curve of $\tan \varphi$ vs. $c /(\gamma H)$, its minimum FOS would be greater than 1.000 and the slope is in a stable state, otherwise, the slope state is unstable with the minimum FOS being less than


Fig. 7 Stability charts for optimizing slope drainage design with slope angle $\beta=45^{\circ}$

Table 3 Comparison of the calculated results

| Case | Slopeparameters |  | Soil parameters |  |  |  |  | Drainage parameters |  | Designed minimum FOS ofslope | Calculated minimum FOS of slope |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $H(\mathrm{~m})$ | $\beta\left({ }^{\circ}\right)$ | $\gamma\left(\mathrm{kN} / \mathrm{m}^{3}\right)$ |  |  |  |  | $l_{w} \tan \beta / H$ | $h_{w} / H$ |  |  |
| 1 | 10 | 45 | 17.8 | 20 | 10 | 18 | 20 | 0.0 | 0.576 | 1.200 | 1.202 |
| 2 | 10 | 45 | 17.8 | 20 | 10 | 18 | 20 | 0.2 | 0.509 | 1.200 | 1.207 |
| 3 | 10 | 45 | 17.8 | 20 | 10 | 18 | 20 | 0.4 | 0.415 | 1.200 | 1.202 |

1.000. In addition, when a specific slope is given with different strength parameters, its minimum slope FOS can also solved by a simple linear proportional relationship in the stability charts, which is shown in Fig. 6(a) for the example of $\beta=30^{\circ}, h_{w} / H=0.2$, and $\left(l_{w} \tan \beta\right) / H=0.0$ with curve $s_{1} s_{2}$.

Here, Figs. 6, 7, and 8 is suitable for $h_{w} / H=0.2 \sim 0.8$ and $\left(l_{w} \tan \beta\right) / H=0.0 \sim 0.4$, and the application range of slope angle $(\beta)$ is $30^{\circ} \sim 60^{\circ}$.

To verify the applicability of these charts for the optimization of slope drainage design, the slope
parameters are given as slope height $H=10 \mathrm{~m}$ and slope angle $\beta=45^{\circ}$. In the slope, the soil parameters are $\gamma=17.8$ $\mathrm{kN} / \mathrm{m}^{3}, \gamma_{\text {sat }}=20 \mathrm{kN} / \mathrm{m}^{3}, c=18 \mathrm{kPa}$, and $\varphi=20^{\circ} . \gamma_{w}=10$ $\mathrm{kN} / \mathrm{m}^{3}$ for the unit weight of water. Meanwhile, the slope under groundwater seepage is in the unstable state. To improve the slope stability, the drainage pipes are arranged in the slope to descend the groundwater table, and the minimum slope FOS would be required to reach 1.200 . Thus, according to the given slope and soil parameters, the drainage parameters, i.e., the highest position of groundwater table, are inversely calculated for the


Fig. 8 Stability charts for optimizing slope drainage design with slope angle $\beta=60^{\circ}$
corresponding minimum slope FOS of 1.200. The results are listed in Table 3. Furthermore, the slope stability with the designed drainage parameters is re-analyzed, and the results are also presented in Table 3.

Table 3 shows that it is feasible to design the drainage parameters of slope using these charts in Figs. 6-8, and the designed minimum slope FOS is very close to the calculated value. Thus, with the use of the charts in Figs. 68 , it could not only quickly get the slope stability under groundwater seepage but obtain the optimal parameters of drainage design for the slope with specific safety requirements.

## 5. Conclusions

In this work, the slope stability under the groundwater seepage is studied. Here, to obtain easily the infiltration and buoyancy forces of the groundwater on the sliding body, its
seepage model is reasonably simplified. Then, by combining the stress analysis on the slip surface, the LE solutions for the stability of 2D and 3D slopes under groundwater seepage are derived with the general nonlinear M-C strength criterion. By comparing the results on some slope examples, the feasibility of the proposed method is verified. Furthermore, the research shows that:
(1) Under the condition of the low hydraulic gradient of groundwater, a change in the groundwater table would have the small effect on the slope stability.
(2) For the slope under groundwater seepage, its stability is significantly reduced for the reason of the decease in the normal stress on slip surface from the effect of buoyancy force and the increase in the sliding force of slope from the effect the infiltration force, and its failure range would also become larger.
(3) The drawn stability charts can be used to quickly obtain the slope stability under groundwater seepage and also to optimize the drainage parameters for the slope with
specific safety requirements.

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